# Precise FWER Control for Gaussian Related Fields: Riding the SuRF to continuous land - Part 1 

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#### Abstract

The Gaussian Kinematic Formula (GKF) is a powerful and computationally efficient tool to perform statistical inference on random fields and became a well-established tool in the analysis of neuroimaging data. Using realistic error models, recent articles show that GKF based methods for voxelwise inference lead to conservative control of the familywise error rate (FWER) and for cluster-size inference lead to inflated false positive rates. In this series of articles we identify and resolve the main causes of these shortcomings in the traditional usage of the GKF for voxelwise inference. This first part removes the good lattice assumption and allows the data to be non-stationary, yet still assumes the data to be Gaussian. The latter assumption is resolved in part 2, where we also demonstrate that our GKF based methodology is non-conservative under realistic error models.


## 1 Introduction

In experiments in neuroscience using functional Magnetic Resonance Imaging (fMRI) of the brain, data consists of 3D-images representing the time dynamics of the blood-oxygen-level dependence (BOLD). After an extensive pre-processing pipeline, including corrections for head motion and respiratory effects the time-series for each subject is combined into subject level maps (3D-images, typically consisting of coefficients of a linear model) via a first-level analysis step, see e.g., Poldrack and Nichols [2011] and the references therein. These high resolution images contain hundreds of thousands of voxels and allow for the detection of the locations of differences in \%BOLD in the brain across different tasks and subjects. A major practical challenge, however, is the low signal-to-noise ratio which is often increased by convolving each subject level map with a smoothing kernel, typically an isotropic 3D-Gaussian kernel with Full Width at Half Maximum $(\mathrm{FWHM}) \approx 3$ voxels [Woo et al., 2014]. A standard statistical analysis to identify areas of activation is the mass-univariate approach. It combines hypotheses tests at each voxel of a test statistic image $\tilde{T}$ obtained from the smoothed subject level maps with a multiple testing procedure. A practical solution to the multiple testing of the thousands of voxels in the brain in terms of controlling the family-wise error rate (FWER) was pioneered in Worsley et al. [1992], Friston et al. [1994], Worsley et al. [1996] and is known as Random Field Theory (RFT) in the neuroimaging community. The key innovation was to use the Gaussian Kinematic Formula (GKF) Adler [1981], Adler and Taylor [2007], Taylor et al. [2006] to approximate the probability that the maximum of a Gaussian random field over a compact
manifold $\mathcal{M} \subset \mathbb{R}^{D}, D \in \mathbb{N}$, with $C^{3}$ sample paths exceeds a given threshold. The accuracy of this approximation has been theoretically validated in [Taylor et al., 2005]. We will call this methodology, where $\mathcal{M}$ is either the brain or the surface of the brain, traditional RFT. Under the assumption of stationarity of the data, traditional RFT became a standard tool in the analysis of fMRI data, Eklund et al. [2016], because of its low computational costs. In the late 1990s it even became the core inference methodology in standard software such as Statistical Parametric Mapping, [Friston et al.] and is still the backbone of peak and cluster-size inference Chumbley and Friston [2009], Friston et al. [1994]. However, in the early 2000s it was shown in Monte-Carlo-simulations that RFT for voxel-wise inference is quite conservative [Nichols and Hayasaka, 2003]. Therefore it has been superseded by time costly, yet accurate permutation approaches, among others, [Nichols and Holmes, 2002, Winkler et al., 2016a,b]. More recently, in the seminal article Eklund et al. [2016], it was demonstrated using resting state data combined with fake task designs that voxelwise RFT inference is conservative (see their Figure 1). They also showed that cluster-size inference based on RFT can have inflated false positive rates and argued that a lack of stationarity and sufficient smoothness of the data are among the causes for these findings. Although these play a role in the conservativeness of voxel-wise inference using RFT the main causes are a mismatch between the probabilistic model and the data and the non-Gaussianity of the data. Here we will present a solution to the conservativeness under the assumption that the data is Gaussian, while part 2 of this paper series Davenport et al. will generalize this approach to non-Gaussian data and validate its performance using a large resting-state dataset from the UK Biobank [Alfaro-Almagro et al., 2018].

For example in [Nichols and Hayasaka, 2003] is has been argued that the lack of sufficient smoothness causes the conservative voxel-wise inference in traditional RFT. Instead we will show that is a relic of a mismatch between probabilistic theory and data. Traditional RFT assumes that an underlying continuous random field is approximated well by the random field observed on the voxel lattice. We call this the good lattice assumption which has been expressed in the following statement from [Worsley et al., 1996]: "[...] the search region was regarded as a region with a smooth boundary defined at every point in 3-D. In practice only voxel data are available, and this will be regarded as a continuous image sampled on a lattice of equally spaced points. Thus a voxel is treated as a point in 3-D with zero volume, although it is often displayed on computer screens and in publications as a volume centered at that point."

This mixes two concepts: the atoms of a probabilistic model and the data. (Raw) data is the information about nature extracted from measurement devices which without exception is currently discrete. In contrast, the atoms (or data objects ${ }^{1}$ ) of probabilistic models used to describe data and perform statistical inference can be functions, curves, images, shapes, trees or other complex mathematical objects [Wang and Marron, 2007, Marron and Alonso, 2014]. In our neuroimaging example, data consists of the subject level maps obtained from the first level model observed on a discrete set $\mathcal{V}$ of voxels which "belong" to the brain. In traditional RFT also the test statistic $\tilde{T}$ used for inference is a random field over $\mathcal{V}$ obtained from the smoothed data, while the atoms of RFT are random fields with twice differentiable sample paths over the brain, i.e., a Whitney stratified (WS) manifold $\mathcal{M}$. Historically, this mismatch has been argued away by the good lattice assumption, which mathematically can be made more precise by assuming $\max _{v \in \mathcal{V}} \tilde{T}(v) \approx \max _{s \in S} \tilde{T}(s)$. This is only reasonable for a high level of applied smoothing, [Kie, 1999, Nichols and Hayasaka, 2003]. Otherwise the FWER control is conservativeness because

$$
\begin{equation*}
\mathrm{P}\left(\max _{v \in \mathcal{V}} \tilde{T}(v)>u_{\alpha}\right) \leq \mathrm{P}\left(\max _{x \in \mathcal{M}} \tilde{T}(x)>u_{\alpha}\right) \approx \alpha \tag{1}
\end{equation*}
$$

[^0]if $u_{\alpha} \in \mathbb{R}$ is the $\alpha$ level threshold obtained from the GKF. To solve this, Worsley [2005] and Taylor et al. [2007] proposed voxel based methods, i.e., estimating the l.h.s. of (1), to control the FWER under the (unrealistic) assumption that the data (the unsmoothed subject level maps) consist of signal plus Gaussian white noise.

We resolve the inconsistency between the atoms and the data and thereby the conservativeness differently. Our key observation is that the smoothing, which is used in pre-processing to increase the signal-to-noise ratio, transforms the data into functions over a WS manifold $\mathcal{M} \supset \mathcal{V}$ where $\mathcal{M}$ depends on the smoother and usually can be specified. Thus, the smoothed subject level maps and $\tilde{T}$ are functions over $\mathcal{M}$ and therefore atoms of RFT. This view allows to estimate quantities required in the GKF such as the Lipschitz Killing Curvatures (LKCs) from the whole functions and not only from discretizations on $\mathcal{V}$ and therefore results in better estimates of $u_{\alpha}$.

More concretely, assume the data $X$ is a realization of a random field over a finite, discrete set $\mathcal{V} \subset \mathbb{R}^{D}, D \in \mathbb{N}$. Using a kernel function $K: \mathcal{V} \times \mathcal{M} \rightarrow \mathbb{R}$ we transform it into functions over a compact, $d$-dimensional WS manifold $\mathcal{M}$, the transformed data, which is

$$
\begin{equation*}
\tilde{X}(s)=\sum_{v \in \mathcal{V}} X(v) K(v, s) . \tag{2}
\end{equation*}
$$

As the latter can be evaluated at all $s \in \mathcal{M}$ we call $\tilde{X}$ a SUper Resolution Field (SuRF). Since SuRFs are a class of random fields, for which many path properties (e.g., differentiability) are inherited directly from the kernel function $K$, we are able to show in Theorem 1 that the GKF holds for SuRFs under mild conditions on $K$ and the data process $X$.

In order to perform inference using the GKF we still must specify the domain $\mathcal{M}$. For many applications, e.g., neuroimaging, it is sensible to ensure that $\mathcal{M} \supset \mathcal{V}$. We propose to use voxel manifolds, which are introduced in Section 3.3. They are composed of the union of $D$-dimensional hyperrectangles which are centered at the elements of $\mathcal{V}$, are aligned with the standard coordinate axis and have as edge lengths $\min \left\{\left|v_{d}-w_{d}\right| \mid v, w \in \mathcal{V}: v_{d} \neq w_{d}\right\}$ in the $d$ th coordinate direction. This is a natural way of defining a WS manifold from $\mathcal{V}$ and was used in Worsley et al. [1996] to approximate the volumes of their underlying unknown search region. For us the main benefit is that the boundary of $\mathcal{M}_{\mathcal{V}}$ consists of the union of lower dimensional hyperrectangles parallel to the coordinate axes. This simplifies the computation the Riemannian metric, induced on $\mathcal{M}_{\mathcal{V}}$ by $\tilde{X}$, and other geometric quantities such as the LKCs, see Section 3.4 and Proposition 7. As $\mathcal{M}_{\mathcal{V}}$ is a WS manifold of dimension greater than zero the map $X \mapsto \tilde{X}$ transforms the data $X$ into atoms of RFT if the functions $s \mapsto K(v, s)$ for all $v \in \mathcal{V}$ are sufficiently regular. Moreover, as $\mathcal{M}_{\mathcal{V}}$ is also the domain of the test statistic $\tilde{T}$ (which depends on transformed data), local maxima and excursions of $\tilde{T}$ above a given threshold can be obtained using numerical optimizers. The latter removes the conservativeness of traditional RFT caused by (1).

In current software packages the LKCs are estimated from the random fields evaluated on $\mathcal{V}$ under stationarity (Forman et al. [1995], [Kie, 1999], Worsley et al. [1996]). These approaches use discrete approximations of the derivatives of $\tilde{X}$ instead of the true derivatives of $\tilde{X}$. The warping estimator [Tay, 2007], and the Hermite projection estimator (HPE) [Telschow et al., 2023] remove the restrictive stationarity assumption, but still rely on discrete approximations and are difficult to implement. For estimators such as the HPE, evaluating the SuRFs at a high resolution reduces the effect of the discretization. This comes at the price of long computation times as we show in our simulations, compare Table 1. In Section 3.5, we propose SuRF based LKC estimators, which use the explicit formulas for the LKCs up to $D=3$ and neither rely
on stationarity nor discrete approximations of derivatives. Moreover, they are computational efficient as they require only a minimal resolution increase to obtain accurate estimates of the LKCs.

As a further contribution we shed some light on effect localization in the data from methods controlling FWER in the strong sense for SuRFs. Any statistical inference methodology relying on a test statistic derived from fields of the form (2) can only localize effects up to the support of the kernel functions $K$. Although important, the exact nature of the FWER control provided has, to the best of our knowledge, not been stated clearly in the literature. Even widely cited articles such as Nichols and Hayasaka [2003] refer only to methods having FWER control in the strong sense, yet do not clearly explain that the strong control is only with respect to the smoothed signal (an atom) and not on the level of data.

The article is structured as follows. In Section 2 we define the notation. Section 3 studies theoretical properties of SuRFs. Section 3.6, explains in detail how voxelwise inference based on the GKF for SuRFs is carried out and provides an explanation of the type of FWER control that this implies. In Section 4.2 the results of the SuRF LKC estimates are reported under different models and are compared to other published LKC estimators. Simulations which demonstrate that the SuRF methodology avoids the conservativeness of traditional voxelwise RFT inference can be found in Section 4.3. In Section 5 we discuss our findings and review other potential applications of our SuRF methodology. A Matlab implementation of the SuRF methodology is available in the RFTtoolbox [Davenport and Telschow, 2023]. Scripts reproducing our simulation results are available at https://github.com/ftelschow/ConvolutionFieldsTheory.

## 2 Notation and Definitions

In this section, we establish notation used in the article. We assume that $v \in \mathbb{R}^{D}$ is a column vector, i.e., we identify $\mathbb{R}^{D}$ with $\mathbb{R}^{D \times 1}$. With $\mathcal{M}$ we denote a $D$-dimensional, compact $\mathcal{C}^{2}$ -Whitney-stratified (WS) manifold isometrically embedded into a $D$-dimensional manifold $\overline{\mathcal{M}}$ without boundary. Recall that a WS manifold of dimension $D$ is a space $\mathcal{M}=\bigcup_{d=1}^{D} \partial_{d} \mathcal{M}$ decomposed into strata where the stratum $\partial_{d} \mathcal{M} \subset \overline{\mathcal{M}}$ is a manifold of dimension $d$ and all the strata are disjoint, i.e., $\partial_{d} \mathcal{M} \cap \partial_{d^{\prime}} \mathcal{M}=\emptyset$ for all $d=d^{\prime}$, compare [Adler and Taylor, 2007, Chapter 8] for more details. A $\mathcal{C}^{2}$-chart (which gives local $\mathcal{C}^{2}$-coordinates) around $s \in \overline{\mathcal{M}}$ is given by a tuple $(\bar{U}, \bar{\phi})$ where $\bar{U} \subset \overline{\mathcal{M}}$ open, $s \in \bar{U}$ and $\bar{\phi} \in \mathcal{C}^{2}(\bar{U}, \bar{V})$ is a diffeomorphism onto an open set $\bar{V} \subset \mathbb{R}^{D}$. By the compactness of $\mathcal{M}$ there exists a set of finitely many charts $\left(\bar{U}_{\alpha}, \bar{\phi}_{\alpha}\right)_{\alpha \in\{1, \ldots, P\}}$ of $\mathcal{M}, P \in \mathbb{N}$, such that $\mathcal{M} \subset \bigcup_{\alpha=1}^{P} \bar{U}_{\alpha}$. This union is the only relevant part of $\overline{\mathcal{M}}$ since we are only interested in properties of $\mathcal{M}$. The surrounding manifold $\overline{\mathcal{M}}$ is only introduced for technical requirements in the formulation of the GKF.

We denote with $f_{\alpha}=f \circ \phi_{\alpha}^{-1}$ the coordinate representation of $f$ in the chart $U_{\alpha}$ and with $\nabla f_{\alpha}$ the gradient of $f_{\alpha}$, i.e., $\nabla f_{\alpha}(x)=\left(\frac{\partial f_{\alpha}}{\partial x_{1}}(x), \ldots, \frac{\partial f_{\alpha}}{\partial x_{D}}(x)\right) \in \mathbb{R}^{1 \times D}$ for $x \in V_{\alpha}$, and with $\nabla^{2} f_{\alpha}$ the Hessian of $f_{\alpha}$, i.e., $\nabla^{2} f_{\alpha}(x) \in \mathbb{R}^{D \times D}$ is the matrix with $d$ - $d^{\prime}$ th entry $\frac{\partial^{2} f_{\alpha}}{\partial x_{d} \partial x_{d^{\prime}}}(x)$ for $x \in V_{\alpha}$. If the gradient $\nabla$ or the Hessian $\nabla^{2}$ is applied to a function with two arguments, then it is always assumed to be with respect to the first argument. We will use $s, s^{\prime}$ for points in $\mathcal{M}$ or $\overline{\mathcal{M}}$ and $x, y$ for points in local coordinates. For simplicity in notation, given $h \in C^{1}\left(\mathbb{R}^{D}\right)$ and a multi-index $\beta \in \mathbb{N}^{d}, d \leq D$, we write

$$
\begin{equation*}
\partial_{\beta} h(x)=\frac{\partial^{d} h(x)}{\partial x_{\beta_{1}}, \ldots, \partial x_{\beta_{d}}}(x), \quad x \in \mathbb{R}^{D} . \tag{3}
\end{equation*}
$$

If $h \in C^{1}\left(\mathbb{R}^{D} \times \mathbb{R}^{D}\right)$, we sometimes write $\left(\partial_{\beta}^{x} h\right)(x, y)$ for $\left(\partial_{\beta} h(\cdot, y)\right)(x)$ evaluated at $(x, y) \in$ $\mathbb{R}^{D} \times \mathbb{R}^{D}$. Note that " $\partial$ " also appears in our notation for the strata of a WS manifold. For a symmetric matrix $A \in \mathbb{R}^{D \times D}, D \in \mathbb{N}$, we define its half-vectorization as $\mathbb{V}(A)=$ $\left(A_{11}, \ldots, A_{D 1}, A_{22}, \ldots, A_{D 2}, \ldots, A_{D-1 D-1}, A_{D D-1}, A_{D D}\right)$ and the set $a \cdot \mathbb{Z}^{D}=\left\{x \in \mathbb{R}^{D} \mid x=\right.$ $\left.\left(a_{1} z_{1}, \ldots, a_{D} z_{N}\right), z \in \mathbb{Z}^{D}\right\}$ for $D \in \mathbb{N}$ and any vector $a \in \mathbb{R}^{D}$ which has positive entries. We also assume throughout the article that $\mathcal{V}$ is a discrete set.

## 3 Theory

In this section we define Super-Resolution Fields (SuRFs), introduce some of their basic properties and show that they satisfy the Gaussian kinematic formula (GKF) under mild conditions. Moreover, we derive computable estimators of the Lipschitz Killing curvatures (LKCs) of a SuRF defined over a voxel manifolds even if the data is non-stationary.

### 3.1 Super-resolution Fields

Definition 1. We call a map $K: \overline{\mathcal{M}} \times \mathcal{V} \rightarrow \mathbb{R}$ a kernel. We say it is continuous/differentiable, if $s \mapsto K(s, v)$ is continuous/differentiable for all $v \in \mathcal{V}$.

Remark 1. This kernel definition is broader than the standard one in statistics, where it is typically a function $k: \overline{\mathcal{M}} \rightarrow \mathbb{R}$ with a normalization property. The function $k$ can be viewed as a kernel from $\overline{\mathcal{M}} \times \overline{\mathcal{M}} \rightarrow \mathbb{R}$ by setting $K\left(s, s^{\prime}\right)=k\left(s-s^{\prime}\right)$.

Definition 2. Let $X$ be a $\mathbb{R}$-valued random field on $\mathcal{V}$ with covariance function $\mathfrak{c}(u, v)=$ $\operatorname{Cov}[X(u), X(v)], u, v \in \mathcal{V}$. For a kernel $K: \overline{\mathcal{M}} \times \mathcal{V} \rightarrow \mathbb{R}$ the random field over $\overline{\mathcal{M}}$,

$$
\begin{equation*}
\tilde{X}(s)=\sum_{v \in \mathcal{V}} K(s, v) X(v), \quad s \in \overline{\mathcal{M}}, \tag{4}
\end{equation*}
$$

is termed a Super-Resolution Field (SuRF) linked to $K$ and $\mathcal{V}$. For brevity, we will refer to it as $(\tilde{X}, X, K, \mathcal{V})$ or simply $\tilde{X}$ when the context is evident regarding $X, K$, and $\mathcal{V}$.

Given a $\operatorname{SuRF}(\tilde{X}, X, K, \mathcal{V})$ we call $(\tilde{X}, X, \tilde{K}, \mathcal{V})$ a normalized $S u R F$ if

$$
\tilde{K}: \overline{\mathcal{M}} \times \mathcal{V} \rightarrow \mathbb{R}, \quad(s, v) \mapsto \frac{K(s, v)}{\sqrt{\sum_{u \in \mathcal{V}} \sum_{v \in \mathcal{V}} K(s, u) K(s, v) \mathfrak{c}(u, v)}}
$$

Remark 2. Let $\tilde{X}$ be a SuRF linked to $K$ and $\mathbb{E}[X(v)]<\infty$ for all $v \in \mathcal{V}$. Then by linearity

$$
\begin{equation*}
\mathbb{E}[\tilde{X}(s)]=\sum_{v \in \mathcal{V}} K(s, v) \mathbb{E}[X(v)], \quad s \in \overline{\mathcal{M}} \tag{5}
\end{equation*}
$$

and if additionally $\mathbb{E}\left[X(v)^{2}\right]<\infty$ for all $v \in \mathcal{V}$ then

$$
\begin{equation*}
\operatorname{Cov}\left[\tilde{X}(s), \tilde{X}\left(s^{\prime}\right)\right]=\sum_{u \in \mathcal{V}} \sum_{v \in \mathcal{V}} K(s, u) K\left(s^{\prime}, v\right) \mathfrak{c}(u, v)<\infty, \quad s, s^{\prime} \in \overline{\mathcal{M}} . \tag{6}
\end{equation*}
$$

The last formula explains the definition of the normalized SuRF, since it shows that any normalized SuRF satisfies $\operatorname{Var}[\tilde{X}(s)]=1$ for all $s \in \overline{\mathcal{M}}$.

Remark 3. Equation (6) is similar to an inner product, where $(\mathfrak{c}(u, v))_{u, v \in \mathcal{V}}$ is the representing "matrix". As such, we introduce, for all $s, s^{\prime} \in \overline{\mathcal{M}}$, the simplifying abbreviations

$$
\begin{equation*}
\left\langle K_{s}, K_{s^{\prime}}\right\rangle=\sum_{u, v \in \mathcal{V}} K(s, u) K\left(s^{\prime}, v\right) \mathfrak{c}(u, v), \quad\left\|K_{s}\right\|^{2}=\sum_{u, v \in \mathcal{V}} K(s, u) K(s, v) \mathfrak{c}(u, v) . \tag{7}
\end{equation*}
$$

Remark 4. If $\mathcal{V} \subseteq \overline{\mathcal{M}}$ and $k: \overline{\mathcal{M}} \rightarrow \mathbb{R}$, then we call the SuRF obtained from the kernel given in Remark 1 a convolution field. Convolution fields appear naturally in many applications since convolving the observed data with a smoothing kernel is often a preprocessing step in signal processing or neuroimaging to improve the signal to noise ratio [Turin, 1960, Worsley et al., 2002].

SuRFs are random fields with nice properties as most path properties are directly inherited from the kernel $K$. This is theoretically advantageous, as one can often establish the assumptions of results like the GKF by imposing conditions like differentiability on the kernel $K$. The following proposition is self-evident but provided here for convenience.
Proposition 1. Let $(\tilde{X}, X, K, \mathcal{V})$ be a SuRF and $k \geq 0$. If $K(\cdot, v) \in \mathcal{C}^{k}(\overline{\mathcal{M}})$ for all $v \in \mathcal{V}$, then $\tilde{X}$ has sample paths of class $\mathcal{C}^{k}$.

### 3.2 Gaussian Kinematic Formula for SuRFs

To demonstrate the benefit of thinking in terms of SuRFs we pose assumptions on a kernel $K$ and the discrete field $X$ such that the corresponding normalized SuRF $\tilde{X}$ satisfies the assumptions of the GKF [Adler and Taylor, 2007, Theorem 12.4.1, 12.4.2]. Thus, we first state the assumptions on a random field $f$ defined over $\overline{\mathcal{M}}$ such that the GKF over $\mathcal{M}$ holds. Recall that $f_{\alpha}$ is the representation of $f$ in the chart $\left(\bar{U}_{\alpha}, \bar{\phi}_{\alpha}\right)$ as introduced in Section 2. In this notation the GKF holds, if for all $\alpha \in\{1, \ldots, P\}$ and $\left(\bar{U}_{\alpha}, \bar{\phi}_{\alpha}\right)$ from the atlas of $\overline{\mathcal{M}}$ we have that
(G1) $f$ is a zero-mean, unit-variance and Gaussian on $\overline{\mathcal{M}}$ with a.s. $\mathcal{C}^{2}$-sample paths.
(G2) $\left(\nabla f_{\alpha}(x), \mathbb{V}\left(\nabla^{2} f_{\alpha}(x)\right)\right)$ is non-degenerate for all $x \in \bar{\phi}_{\alpha}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$.
(G3) There exist constants $\kappa, \gamma, \epsilon>0$ such that for each $d, d^{\prime} \in\{1, \ldots, D\}$,

$$
\mathbb{E}\left[\left(\partial_{d d^{\prime}} f_{\alpha}(x)-\partial_{d d^{\prime}} f_{\alpha}(y)\right)^{2}\right] \leq \kappa|\log \|x-y\||^{-1-\gamma}
$$

for all $x, y \in \bar{\phi}_{\alpha}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ for which $|x-y|<\epsilon$.
Remark 5. Conditions (G1)-(G3) imply that the sample paths of $f$ are almost surely Morse functions, compare [Adler and Taylor, 2007, Corollary 11.3.2.]. Moreover, by Lemma 1 from Davenport and Telschow [2022] these conditions do not depend on the particular choice of the $\mathcal{C}^{3}$ charts $\left(U_{\alpha}, \phi_{\alpha}\right), \alpha \in\{1, \ldots, P\}$, but rather hold for all $\mathcal{C}^{3}$ charts $(V, \varphi)$ of $\overline{\mathcal{M}}$.
Applying Proposition 1, (G1) holds for any normalized SuRF derived from a zero-mean Gaussian random field $X$ on $\mathcal{V}$ and a twice continuously differentiable kernel $K$. The next two propositions establish that Condition (G3) is satisfied for a SuRF with a $\mathcal{C}^{3}$-kernel.

Proposition 2. Let $\gamma \in(0,1]$ and for all $v \in \mathcal{V}$ let $K(\cdot, v)$ be $\gamma$-Hölder continuous with Hölder constants bounded above by $A>0$ and $\mathbb{E}\left[X(v)^{p}\right]<\infty, p \in[1, \infty)$. Then $\tilde{X}$ has almost surely $\mathcal{L}^{p}$-Hölder continuous paths, i.e.,

$$
\begin{equation*}
\left|\tilde{X}_{\alpha}(x)-\tilde{X}_{\alpha}(y)\right| \leq L\|x-y\|^{\gamma} \tag{8}
\end{equation*}
$$

for the charts $\left(\bar{U}_{\alpha}, \bar{\phi}_{\alpha}\right), \alpha \in\{1, \ldots, P\}$, in the atlas of $\overline{\mathcal{M}}$ covering $\mathcal{M}$, all $x, y \in \bar{\phi}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ and some random variable $L$ with finite $p$-th moment.

Proposition 3. Let $K(\cdot, v) \in \mathcal{C}^{1}(\overline{\mathcal{M}})$ and $\mathbb{E}\left[X(v)^{2}\right]<\infty$ for all $\nu \in \mathcal{V}$. Then there exists a constant $\kappa>0$ for $\alpha \in\{1, \ldots, P\}$ such that

$$
\begin{equation*}
\mathbb{E}\left[\left(\tilde{X}_{\alpha}(x)-\tilde{X}_{\alpha}(y)\right)^{2}\right] \leq \kappa|\log \|x-y\||^{-2} \tag{9}
\end{equation*}
$$

for all $x, y \in \bar{\phi}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ such that $0<\|x-y\|<1$.
Remark 6. Proposition 3 implies condition (G3) for a $\mathcal{C}^{3}$-kernel $K$ because each second order partial derivative of a normalized SuRF is a SuRF with continuous sample paths.

Definition 3. Given a set $\mathcal{W} \subset \mathbb{R}^{D}$, we say that functions $f_{1}, \ldots, f_{J}: \mathcal{W} \rightarrow \mathbb{R}$ are $\mathcal{W}$-linearly independent if given constants $a_{1}, \ldots, a_{J} \in \mathbb{R}$, the relation

$$
\sum_{j=1}^{J} a_{j} f_{j}(w)=0
$$

holding for all $w \in \mathcal{W}$ implies that $a_{j}=0$ for all $j \in\{1, \ldots, J\}$.
The following propositions link the linear independence of functions $K_{\alpha}(x, \cdot)$ for $x \in \bar{\phi}_{\alpha}\left(\bar{U}_{\alpha} \cap\right.$ $\mathcal{M}$ ) with Condition (G2) for SuRFs. We demonstrate that the needed linear independence condition for achieving (G2) is satisfied, for example, by Gaussian kernels.

Proposition 4. Let $(\tilde{X}, X, K, \mathcal{V})$ be a SuRF with $K(\cdot, v) \in C^{2}(\overline{\mathcal{M}})$ for all $v \in \mathcal{V}$ and $\tilde{Z}$ be the corresponding normalized SuRF. For $\alpha \in\{1, \ldots, P\}$ and $x \in \bar{\phi}_{\alpha}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ define $K_{\alpha}(x, \cdot): \mathcal{V} \rightarrow$ $\mathbb{R}$ by $v \mapsto K\left(\bar{\phi}_{\alpha}^{-1}(x), v\right)$ and set $\mathcal{V}_{x}=\left\{v \in \mathcal{V}: K_{\alpha}(x, v) \neq 0\right\}$. Assume that $K_{\alpha}(x, \cdot), \partial_{d}^{x} K_{\alpha}(x, \cdot)$ and $\partial_{d^{\prime} d^{\prime \prime}}^{x} K_{\alpha}(x, \cdot)$ for $d, d^{\prime}, d^{\prime \prime} \in\{1, \ldots, D\}$ with $1 \leq d^{\prime} \leq d^{\prime \prime} \leq D$ are $\mathcal{V}_{x}$-linearly independent and the random vector $\left(X(v): v \in \mathcal{V}_{x}\right)$ is non-degenerate. Then

$$
\left(\tilde{X}_{\alpha}(x), \nabla \tilde{X}_{\alpha}(x), \mathbb{V}\left(\nabla^{2} \tilde{X}_{\alpha}(x)\right)\right) \text { and }\left(\tilde{Z}_{\alpha}(x), \nabla \tilde{Z}_{\alpha}(x), \mathbb{V}\left(\nabla^{2} \tilde{Z}_{\alpha}(x)\right)\right)
$$

are non-degenerate Gaussian random vectors.
Proposition 5. Let $\tilde{D}=D+1+D(D+1) / 2$ and $\alpha \in\{1 \ldots, P\}$. Taking the gradient and Hessian w.r.t. $x \in \bar{\phi}_{\alpha}\left(\bar{U}_{\alpha} \cap \mathcal{M}\right)$ we define the vector valued functions

$$
\begin{equation*}
\mathbf{K}_{\alpha, x}(v)=\left(K_{\alpha}(x, v), \nabla K_{\alpha}(x, v), \mathbb{V}\left(\nabla^{2} K_{\alpha}(x, v)\right)\right) \tag{10}
\end{equation*}
$$

indexed by $x \in \bar{\phi}_{\alpha}\left(\bar{U}_{\alpha} \cap \mathcal{M}\right)$. If for each $x \in \bar{\phi}_{\alpha}\left(\bar{U}_{\alpha} \cap \mathcal{M}\right)$ there exist $v_{1}, \ldots, v_{\tilde{D}} \in \mathcal{V}_{x}$ such that $\mathbf{K}_{\alpha, x}\left(v_{1}\right), \ldots, \mathbf{K}_{\alpha, x}\left(v_{\tilde{D}}\right)$ are linearly independent, then $K_{\alpha}(x, \cdot), \partial_{d}^{x} K_{\alpha}(x, \cdot)$ and $\partial_{d^{\prime} d^{\prime \prime}}^{x} K_{\alpha}(x, \cdot)$ for $d, d^{\prime}, d^{\prime \prime} \in\{1, \ldots, D\}$ with $1 \leq d^{\prime} \leq d^{\prime \prime} \leq D$ are $\mathcal{V}_{x}$ linearly independent.

Remark 7. The results in Propositions 4 and 5 are stronger than (G2) as we want to emphasize that SuRFs often satisfy the assumptions of the expectation Metatheorem 11.2.1 from Adler and Taylor [2007], see Corollary 11.2.2 for the Gaussian version. However, our proofs show that (G2) follows already from the assumption that $\partial_{d}^{x} K(x, \cdot)$ and $\partial_{d^{\prime} d^{\prime \prime}}^{x} K(x, \cdot)$ for $d \in\{1, \ldots, D\}$ and $1 \leq d^{\prime} \leq d^{\prime \prime} \leq D$ is $\mathcal{V}_{x}$-linearly independent for all $x \in \bar{\phi}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ and $\alpha \in\{1, \ldots, P\}$.

Proposition 6. Let $K(s, v)=e^{-(s-v)^{T} \Omega(s-v) / 2}$, for some positive definite matrix $\Omega \in \mathbb{R}^{D \times D}$ be the $D$-dimensional Gaussian kernel. Assume that $\mathcal{V}$ is a $D$-dimensional lattice which contains an element $v$ such that

$$
\left\{v+\sum_{d=1}^{D} \lambda_{d} \gamma_{d} e_{d}: \gamma_{d} \in\{-1,0,1\}\right\} \subset \mathcal{V}
$$

where $\left(e_{d}\right)_{1 \leq d \leq D}$ is the standard basis and $\lambda \in \mathbb{R}_{>0}^{D}$. Then $K(s, \cdot), \partial_{d}^{s} K(s, \cdot)$ and $\partial_{d^{\prime} d^{\prime \prime}}^{s} K(s, \cdot)$ for $d \in\{1, \ldots, D\}$ and $1 \leq d^{\prime} \leq d^{\prime \prime} \leq D$ are $\mathcal{V}_{s}$-linearly independent for each $s \in \mathbb{R}^{D}$.

Having established the conditions above we can now proof the GKF for Gaussian related fields obtained from Gaussian SuRFs. It is a corollary to Theorem 12.4.2 from Adler and Taylor [2007]. To state it we define $\mathcal{A}_{u}(f)=\{s \in \mathcal{M}: f(s) \geq u\}$ to be the excursion set of a random field $f$ above the threshold $u$ on $\mathcal{M}$ and write $\chi_{f}(u)$ to denote the EC of the excursion set $\mathcal{A}_{u}(f)$.

Theorem 1. Let $\left(\tilde{X}_{1}, X_{1}, K, \mathcal{V}\right), \ldots,\left(\tilde{X}_{N}, X_{N}, K, \mathcal{V}\right) \sim(\tilde{X}, X, K, \mathcal{V})$ be i.i.d. SuRFs and $F \in$ $\mathcal{C}^{2}(\overline{\mathcal{M}})$. Assume that $X$ is a Gaussian field on $\mathcal{V}$ with covariance function $\mathfrak{c}$ and that for all $v \in \mathcal{V}$ it holds that $K(\cdot, v) \in C^{3}(\overline{\mathcal{M}})$ and $\mathfrak{c}(v, v)>0$. Furthermore, for all $\alpha \in\{1, \ldots, P\}$ and $x \in \bar{\phi}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ assume that the random vector $\left(X(v): v \in \mathcal{V}_{x}\right)$ is non-degenerate for the $\mathcal{V}_{x}$ defined in Proposition 4 and that $\partial_{d}^{x} K_{\alpha}(x, \cdot)$ and $\partial_{d^{\prime} d^{\prime \prime}}^{x} K_{\alpha}(x, \cdot)$ for $d \in\{1, \ldots, D\}$ and $1 \leq d^{\prime} \leq d^{\prime \prime} \leq D$ are $\mathcal{V}_{x}$-linearly independent. Define a random field $T$ such that $T(s)=$ $F\left(\tilde{X}_{1}(s) /\left\|K_{s}\right\|, \ldots, \tilde{X}_{N}(s) /\left\|K_{s}\right\|\right)$ for all $s \in \overline{\mathcal{M}}$. Then

$$
\begin{equation*}
\mathbb{E}\left[\chi_{T}(u)\right]=\sum_{d=0}^{D} \mathcal{L}_{d} \rho_{d}^{T}(u), \quad u \in \mathbb{R} \tag{11}
\end{equation*}
$$

where $\mathcal{L}_{0}, \ldots, \mathcal{L}_{D} \in \mathbb{R}$ are the LKCs of $\overline{\mathcal{M}}$ endowed with the induced Riemannian metric from $\tilde{X}(s) /\left\|K_{s}\right\|$ and $\rho_{d}^{T}$ 's are functions depending solely on the marginal distribution of $T$.

### 3.3 Voxel Manifolds

Until now we did not discuss what domain $\mathcal{M} \subset \overline{\mathcal{M}}$ we should choose for a Surf derived from a kernel $K$ and a discrete random field $X$ on a finite grid $\mathcal{V}$. In the case that $\overline{\mathcal{M}}$ is a $D$-dimensional submanifold of $\mathbb{R}^{D}$ and $\mathcal{V} \subset \mathbb{R}^{D}$ and $K(\cdot, v): \mathbb{R}^{D} \rightarrow \mathbb{R}$, for each $v \in \mathcal{V}$, we propose to use a practical domain which we call the voxel manifold associated with $\mathcal{V}$.

Definition 4. Suppose that $\mathcal{V} \subset \mathbb{R}^{D}$ is a discrete set and define $\delta \in \mathbb{R}^{D}$ such that its $d$-th component is $\delta_{d}=\min \left\{\left|v_{d}-w_{d}\right|: v, w \in \mathcal{V}, v_{d} \neq w_{d}\right\}$. Moreover, let

$$
\begin{equation*}
\mathcal{B}_{v}(\delta)=\left\{x \in \mathbb{R}^{D}\left|\max _{d \in\{1, \ldots, D\}}\right| x_{d}-v_{d} \mid-\delta_{d} / 2 \leq 0\right\} . \tag{12}
\end{equation*}
$$

Then the voxel manifold associated with $\mathcal{V}$ is the set $\mathcal{M}_{\mathcal{V}}=\bigcup_{v \in \mathcal{V}} \mathcal{B}_{v}(\delta)$.
The advantage of using such a domain for a SuRF is that, for example, estimators of the LKCs can be calculated efficiently, because all $d$-dimensional boundaries lie in hyperplanes parallel to the coordinate axes. In order that a SuRF derived from $K$ and $\mathcal{V}$ satisfies the GKF we need to assume that $\mathcal{B}_{v}(\delta) \subset \operatorname{supp}(K(\cdot, v))$ for all $v \in \mathcal{V}$. Otherwise the $\operatorname{SuRF}$ is zero on parts of $\mathcal{M}_{\mathcal{V}}$ and thus (G2) cannot be true. This condition on the support, however, is usually satisfied for the choice of $K$ as the kernel is typically used to increase the signal-to-noise ratio through averaging observations at different $v \in \mathcal{V}$.

A voxel manifold is a stratified space, for example, for $D=3$ by the three dimensional stratum is the union of all the open cubes $\operatorname{int}\left(\mathcal{B}_{v}(\delta)\right)$ with $v \in \mathcal{V}$, while the two dimensional, the one dimensional and the zero dimensional strata are the unions over all $v \in V$ of all faces, edges and corners of the cubes $\mathcal{B}_{v}(\delta)$ respectively. In fact, a voxel manifold as the union of polyhedra is even a WS manifold [Adler and Taylor, 2007, p.187].

In Riemannian geometry most geometric quantities can be derived from the Riemmanian metric and the Christoffel symbols. Using Theorem 5 and 6 from Appendix B. 1 these quantities for the Riemannian metric induced by a normalized SuRF on $\overline{\mathcal{M}}$ can be written in terms of the inner products introduced in (7).

Proposition 7. For $\overline{\mathcal{M}} \subset \mathbb{R}^{D}$ the Riemannian metric $\boldsymbol{\Lambda}$ induced by a normalized SuRF expressed in the local coordinates induced by $\iota: \overline{\mathcal{M}} \hookrightarrow \mathbb{R}^{D}$ is given by

$$
\begin{equation*}
\Lambda_{d d^{\prime}}(x)=\frac{\left\langle\partial_{d} K_{x}, \partial_{d^{\prime}} K_{x}\right\rangle}{\left\|K_{x}\right\|^{2}}-\frac{\left\langle\partial_{d} K_{x}, K_{x}\right\rangle\left\langle K_{x}, \partial_{d^{\prime}} K_{x}\right\rangle}{\left\|K_{x}\right\|^{4}} \tag{13}
\end{equation*}
$$

and the Christoffel symbols of the first kind are

$$
\begin{align*}
\Gamma_{k d d^{\prime}}(x)= & \frac{\left\langle\partial_{k} \partial_{d} K_{x}, \partial_{d^{\prime}} K_{x}\right\rangle}{\left\|K_{x}\right\|^{2}}-\frac{\left\langle\partial_{k} \partial_{d} K_{x}, K_{x}\right\rangle\left\langle K_{x}, \partial_{d^{\prime}} K_{x}\right\rangle}{\left\|K_{x}\right\|^{4}} \\
& -\frac{\left\langle\partial_{k} K_{x}, K_{x}\right\rangle\left\langle\partial_{d} K_{x}, \partial_{d^{\prime}} K_{x}\right\rangle}{\left\|K_{x}\right\|^{4}}-\frac{\left\langle\partial_{d} K_{x}, K_{x}\right\rangle\left\langle\partial_{k} K_{x}, \partial_{d^{\prime}} K_{x}\right\rangle}{\left\|K_{x}\right\|^{4}}  \tag{14}\\
& +2 \frac{\left\langle\partial_{k} K_{x}, K_{x}\right\rangle\left\langle\partial_{d} K_{x}, K_{x}\right\rangle\left\langle K_{x}, \partial_{d^{\prime}} K_{x}\right\rangle}{\left\|K_{x}\right\|^{6}}
\end{align*}
$$

The advantage of using the voxel manifold domain is that the numerical implementation of geometric quantities is feasible. In particular, we can construct, for any $x \in \mathcal{U}$ and a small open neighborhood $\mathcal{U} \ni x$ an orthonormal basis of $\mathcal{T}_{x} \mathcal{M}_{\mathcal{V}}, x \in \mathcal{M}_{\mathcal{V}} \cap \mathcal{U}$, such that a subset of this basis is an orthonormal frame of $\mathcal{T}_{x} \partial_{d} \mathcal{M}_{\mathcal{V}}, d \in\{1, \ldots, D-1\}$, if $x \in \partial_{d} \mathcal{M}_{\mathcal{V}} \cap \mathcal{U}$. More concretely, for $D=3$ the Gram-Schmidt procedure on the Euclidean basis $e_{1}, e_{2}, e_{3}$ yields the following orthonormal vector fields with respect to the Riemannian metric induced by a SuRF at $x \in \mathcal{M}_{\mathcal{V}}$ and $k, l, m \in\{1,2,3\}$ such that $k<l$ and $\{k, l, m\}=\{1,2,3\}$ :

$$
\begin{gather*}
U_{x}=\Lambda_{k k}^{-1 / 2}(x) e_{k}, \quad V_{x}=\frac{\Lambda_{k l}(x)}{\sqrt{c(x) \Lambda_{k k}(x)}} e_{k}-\sqrt{\frac{\Lambda_{k k}(x)}{c(x)}} e_{l}, \\
N_{x}=\frac{\Lambda^{-1}(x)}{\sqrt{e_{m}^{T} \boldsymbol{\Lambda}^{-1}(x) e_{m}}} e_{m} . \tag{15}
\end{gather*}
$$

Here $c(x)=\operatorname{det}\left(\boldsymbol{\Lambda}^{I}(x)\right)$ for $I=(k, l)$ and $U_{x}, V_{x}$ are a basis of $\mathcal{T}_{x} \mathcal{F}_{I}$, where $\mathcal{F}_{I}$ is the subset of $\partial_{2} \mathcal{M}_{\mathcal{V}}$ such that the coordinates with indices not contained in $I$ are constant, and $N_{x}$ is in the one dimensional vector space orthogonal to $\mathcal{T}_{x} \mathcal{F}_{I}$ with respect to $\boldsymbol{\Lambda}$.

### 3.4 LKCs of Voxel Manifolds

General formulas for the LKCs of WS manifolds can be found in [Adler and Taylor, 2007, Theorem 12.4.2] and formulas for WS manifolds of dimension $D \leq 3$ are given in Appendix B.2. Since voxel manifolds are embedded into $\mathbb{R}^{D}$ and have a simple geometric structure, the
highest two LKCs can be expressed as integrals of (sub-)determinants of the Riemannian metric $\boldsymbol{\Lambda}$, i.e.,

$$
\mathcal{L}_{D}=\sum_{\nu \in \mathcal{V}} \int_{\mathcal{B}_{\nu}(\delta)} \sqrt{\operatorname{det}(\boldsymbol{\Lambda}(x))} d x, \quad \mathcal{L}_{D-1}=\sum_{I:|I|=D-1} \int_{\mathcal{F}_{I}} \sqrt{\operatorname{det}\left(\Lambda^{I}(x)\right)} d x^{I}
$$

Here $I$ is any ordered subset of $\{1, \ldots, D\}, d x^{I}=d x_{I_{1}} \ldots d x_{I_{|I|}}$ and $\boldsymbol{\Lambda}^{I}(x)$ is the submatrix of $\boldsymbol{\Lambda}(x)$ consisting of the columns and rows given by the entries of $I$. Finally, $\mathcal{F}_{I}$ is the subset of $\partial_{|I|} \mathcal{M}_{\mathcal{V}}$ such that the coordinates with indices not contained in $I$ are constant.

The LKCs $\mathcal{L}_{1}, \ldots, \mathcal{L}_{D-2}$ of a $D$-dimensional voxel manifold with $D>2$ are substantially harder to express explicitly, compare [Adler and Taylor, 2007, Theorem 12.4.2]. An expression for $\mathcal{L}_{1}$ of a 3 -dimensional voxel manifold is given in the following proposition.

Theorem 2. Let $\mathcal{M}_{\mathcal{V}}$ be a 3-dimensional voxel manifold. Then

$$
\begin{align*}
\mathcal{L}_{1}= & \frac{1}{2 \pi} \sum_{|I|=1} \int_{\mathcal{F}_{I}} \Theta(x) \sqrt{\operatorname{det}\left(\boldsymbol{\Lambda}^{I}(x)\right)} d x^{I} \\
& +\frac{1}{2 \pi} \sum_{|I|=2} \int_{\mathcal{F}_{I}}\left[\left(U_{I_{1}}(x)^{2}+V_{I_{1}}(x)^{2}\right) N^{T}(x)\left(\begin{array}{c}
\Gamma_{I_{1} I_{1} 1}(x) \\
\Gamma_{I_{1} I_{1} 2}(x) \\
\Gamma_{I_{1} I_{1} 3}(x)
\end{array}\right)\right.  \tag{16}\\
& \left.+\sum_{k=1}^{2} V_{I_{k}}(x) V_{I_{2}}(x) N^{T}(x)\left(\begin{array}{c}
\Gamma_{I_{k} I_{2} 1}(x) \\
\Gamma_{I_{k} I_{2} 2}(x) \\
\Gamma_{I_{k} I_{2} 3}(x)
\end{array}\right)\right] \sqrt{\operatorname{det}\left(\boldsymbol{\Lambda}^{I}(x)\right)} d x^{I} \\
& -\frac{1}{2 \pi} \sum_{v \in \mathcal{V}} \int_{\mathcal{B}_{v}(\delta)} \operatorname{Tr}(R(x)) \sqrt{\operatorname{det}(\boldsymbol{\Lambda}(x))} d x
\end{align*}
$$

Here $\operatorname{Tr}(R(x))$ is the trace of the Riemannian curvature tensor and

$$
\Theta(x)= \begin{cases}\pi-\beta(x), & \text { if } x \text { belongs to a convex edge } \\ -2 \beta(x), & \text { if } x \text { belongs to a double convex edge } \\ \beta(x)-\pi, & \text { if } x \text { belongs to a convex edge }\end{cases}
$$

compare Appendix B.1 and especially (37). Here $\beta(x)$ is defined using the crossproduct $V_{x} \times N_{x}=$ $\left(m_{1}^{I}(x), m_{2}^{I}(x), m_{3}^{I}(x)\right)$ of the elements of the ONB from (15) with $k=I$ by

$$
\beta(x)=\arccos \left(\frac{m_{2}(x) m_{3}(x)}{\sqrt{m_{2}^{2}(x)+m_{1}^{2}(x)} \sqrt{m_{3}^{2}(x)+m_{1}^{2}(x)}}\right)
$$

The different types of edges are visualized in Figure 1.

### 3.5 SuRF Estimator for LKCs

Given a kernel $K: \mathcal{M} \mathcal{V} \times \mathcal{V} \rightarrow \mathbb{R}$ and an i.i.d. sample $X_{1}, \ldots, X_{N}$ of a Gaussian random field $X$ over $\mathcal{V}$, we use formula (13) for the Riemannian metric $\boldsymbol{\Lambda}(x)$ to derive an estimate of the LKCs which does not use discrete derivatives. The idea is to replace variances and covariances by


Figure 1: Visualization of the different types of edges appearing in a voxel manifold. In the left voxel manifold all edges are convex. The edge where the two cubes of the voxel manifold in the center are touching is a double convex edge and the same edge in the voxel manifold on the right is a concave edge since a third cube is added.
their sample counterparts obtained from the sample $\tilde{\mathbb{X}}=\left(\tilde{X}_{1}, \ldots, \tilde{X}_{N}\right)$ of SuRFs. This results in the SuRF-Riemannian metric estimator

$$
\begin{align*}
\hat{\Lambda}_{d d^{\prime}}(x)= & \frac{\operatorname{Cov}\left[\partial_{d} \tilde{\mathbb{X}}(x), \partial_{d^{\prime}} \tilde{\mathbb{X}}(x)\right]}{\operatorname{Var}[\tilde{\mathbb{X}}(x)]} \\
& -\frac{\operatorname{Cov}\left[\partial_{d} \tilde{\mathbb{X}}(x), \tilde{\mathbb{X}}(x)\right] \operatorname{Cov}\left[\tilde{\mathbb{X}}(x), \partial_{d^{\prime}} \tilde{\mathbb{X}}(x)\right]}{\operatorname{Var}[\tilde{\mathbb{X}}(x)]^{2}}, \tag{17}
\end{align*}
$$

where the variances and covariances are sample variances and covariances of the sample $\tilde{\mathbb{X}}$ and its derivatives $\partial_{d} \tilde{\mathbb{X}}(x)$. The latter can be computed from the derivatives of the kernel $K$, compare Proposition 1. Hence, (17) does not rely on numerical derivatives. We denote with $\hat{\Lambda}(x) \in \mathbb{R}^{D \times D}, x \in \mathcal{M}_{\mathcal{V}}$, the matrix with $\left(d, d^{\prime}\right)$-th entry $\hat{\Lambda}_{d d^{\prime}}(x)$. To estimate the LKCs we evaluate $\hat{\boldsymbol{\Lambda}}(x)$ on a grid $\mathcal{M}_{\mathcal{V}}^{(r)} \subset \mathcal{M}_{\mathcal{V}}$ given by

$$
\begin{equation*}
\mathcal{M}_{\mathcal{V}}^{(r)}=\bigcup_{v \in \mathcal{V}} \mathcal{B}_{v}(\delta) \cap\left(v+\frac{\delta}{r+1} \cdot \mathbb{Z}^{D}\right) \tag{18}
\end{equation*}
$$

Here $r \in\left\{2 r^{\prime}+1 \mid r^{\prime} \in \mathbb{N}\right\}$, i.e., $r \geq 1$ and odd, is called the added resolution. The restriction to odd numbers is required to ensure that we sample the boundary of $\mathcal{M}_{\mathcal{V}}$.

The SuRF-LKC estimators with added resolution $r$ for $\mathcal{L}_{D}$ and $\mathcal{L}_{D-1}$ are given by

$$
\begin{aligned}
\hat{\mathcal{L}}_{D}^{(r)} & =\sum_{x \in \mathcal{M}_{\nu}^{(r)}} \sqrt{\operatorname{det}(\hat{\boldsymbol{\Lambda}}(x))} \prod_{d=1}^{D} \frac{\delta_{d}}{r+1} . \\
\hat{\mathcal{L}}_{D-1}^{(r)} & =\sum_{|I|=D-1} \sum_{x \in \mathcal{F}_{I}^{(r)}} \sqrt{\operatorname{det}\left(\hat{\boldsymbol{\Lambda}}^{I}(x)\right)} \prod_{i \in I} \frac{\delta_{i}}{r+1} .
\end{aligned}
$$

Here $\mathcal{F}_{I}^{(r)}=\mathcal{F}_{I} \cap \mathcal{M}_{\mathcal{V}}^{(r)}$. These formulas are easy to implement in software. and as $r$ goes to infinity the numerical error in approximating the integral can be made arbitrary small. Moreover, it is possible to obtain in the same fashion estimators $\hat{\mathcal{L}}_{D-d}^{(r)}, d \in\{1, \ldots D-2\}$ from the LKC formulas in [Adler and Taylor, 2007, Theorem 12.4.2] and our formulas of the geometric quantities induced by a $\operatorname{SuRF}$ on $\mathcal{M}_{\mathcal{V}}$, compare Corollary 7 and Appendix B.2. In practice it is
tedious to implement these estimators and their computation time would be be large, because even for a 3-dimensional voxel manifold we need for the Christoffel symbols 27 convolutions on top of the 9 convolutions required to estimate the Riemannian metric. Computing the Riemannian curvature tensor needs another 36 convolutions. A solution is to approximate the lower LKCs by their locally stationary counterparts, i.e., for $D=3$ only the first integral in Theorem 2 remains:

$$
\begin{equation*}
\hat{\mathcal{L}}_{1}^{(r)}=\sum_{|I|=1} \sum_{v \in \mathcal{F}_{I}^{(r)}} \hat{\Theta}(v) \sqrt{\hat{\boldsymbol{\Lambda}}^{I}(v)} \frac{\delta_{I}}{r+1} . \tag{19}
\end{equation*}
$$

Here $\hat{\Theta}$ is the plug-in estimate of $\Theta$ defined in Theorem 2.
Theorem 3. Let $K$ be the kernel of the SuRF and $d \in\{D-1, D\}$. Assume that $K(\cdot, v) \in \mathcal{C}^{3}$ for all $v \in \mathcal{V}$ and that $\mathbb{E}[X(v)]<\infty$ for all $v \in \mathcal{V}$. Then

$$
\lim _{r \rightarrow \infty} \mathbb{E}\left[\hat{\mathcal{L}}_{d}^{(r)}\right]=\mathbb{E}\left[\lim _{r \rightarrow \infty} \hat{\mathcal{L}}_{d}^{(r)}\right]=\mathcal{L}_{d}
$$

Theorem 4. Let $d \in\{D-1, D\}, K(\cdot, v) \in \mathcal{C}^{3}(\overline{\mathcal{M}})$ and $\mathbb{E}\left[X(v)^{2}\right]<\infty$ for $v \in \mathcal{V}$. Then

$$
\begin{equation*}
\lim _{N \rightarrow \infty} \lim _{r \rightarrow \infty} \hat{\mathcal{L}}_{d}^{(r)}=\lim _{r \rightarrow \infty} \lim _{N \rightarrow \infty} \hat{\mathcal{L}}_{d}^{(r)}=\mathcal{L}_{d} . \tag{20}
\end{equation*}
$$

Remark 8. We expect that a similar result can be derived for the plugin estimator $\mathcal{L}_{1}$ resulting from Theorem 2 where also the $\Gamma_{k d d^{\prime}}$ 's from (14) and the Riemannian curvature, Appendix B.1, are estimated using the corresponding sample covariances. This result could be established along the same lines as the consistency in [Telschow et al., 2023, Section 3], but we leave this for future work since, currently, implementing this estimator seems infeasible.

### 3.6 FWER Control Using SuRFs

In this section we illustrate how the GKF is used to construct a test controlling the familywise error rate (FWER) which dates back to Worsley et al. [1992]. Here we combine this approach with SuRFs in order to improve the power of RFT based voxelwise inference. As an illustration, we consider the problem of detecting areas of non-zero signal $\mu$, given an i.i.d. sample $X_{1}, \ldots, X_{N} \sim X$, in a signal plus noise model $X(v)=\mu(v)+\epsilon(v)$ for $v \in \mathcal{V}$ and $\epsilon$ a zero-mean random field. The same approach can be applied to the linear model and other probabilistic models for which the GKF holds, compare Worsley [1994], Tay [2007], Taylor and Worsley [2008].

Let $\tilde{X}_{1}, \ldots, \tilde{X}_{N}$ denote a sample of SuRFs over $\mathcal{M}_{\mathcal{V}}$ derived from random fields $X_{1}, \ldots, X_{N}$ over $\mathcal{V} .{ }^{2}$ Assume the setting of Theorem 1 and define $F: \mathbb{R}^{N} \rightarrow \mathbb{R}$ by

$$
F\left(a_{1}, \ldots, a_{N}\right)=\frac{1}{\sqrt{N}} \sum_{i=1}^{N} a_{i}\left(\frac{1}{N-1} \sum_{i=1}^{N}\left(a_{i}-\frac{1}{N} \sum_{i=1}^{N} a_{i}\right)^{2}\right)^{-1 / 2}
$$

for $\left(a_{1}, \ldots, a_{N}\right) \in \mathbb{R}^{N}$. Let $\mu(v)=\mathbb{E}[X(v)]$ for $v \in \mathcal{V}$ and for all $s \in \mathcal{M}_{\mathcal{V}}$, let

$$
\tilde{\mu}(s)=\sum_{v \in \mathcal{V}} K(s, v) \mu(v) .
$$

[^1]Due to the invariance of $F$ to scaling, for $s \in \mathcal{M} \mathcal{V}$, we can write

$$
\begin{equation*}
\tilde{T}(s):=F\left(\tilde{X}_{1}(s) /\left\|K_{s}\right\|, \ldots, \tilde{X}_{N}(s) /\left\|K_{s}\right\|\right)=\frac{\sqrt{N} \hat{\mu}_{N}(s)}{\hat{\sigma}_{N}(s)} \tag{21}
\end{equation*}
$$

where $\hat{\mu}_{N}(s)=\frac{1}{N} \sum_{i=1}^{N} \tilde{X}_{i}(s)$ and $\hat{\sigma}_{N}(s)=\left(\frac{1}{N-1} \sum_{i=1}^{N}\left(\tilde{X}_{i}(s)-\hat{\mu}_{N}(s)\right)^{2}\right)^{1 / 2}$.
Our goal is to construct based on $\tilde{T}$ a multiple hypothesis test for the hypotheses

$$
\begin{equation*}
\mathbf{H}_{0, s}: \tilde{\mu}(s) \leq 0 \quad \text { vs. } \quad \mathbf{H}_{1, s}: \tilde{\mu}(s)>0 \quad s \in \mathcal{M}_{\mathcal{V}} \tag{22}
\end{equation*}
$$

which controls the FWER in the strong sense at a significance level $\alpha \in(0,1) .{ }^{3}$ Denote with $\mathcal{H}_{0}=\left\{s \in \mathcal{M}_{\mathcal{V}} \mid \tilde{\mu}(s) \leq 0\right\} \subseteq \mathcal{M}_{\mathcal{V}}$ the set of true null hypotheses and consider the test that rejects $\mathbf{H}_{0, s}, s \in \mathcal{M}_{\mathcal{V}}$, whenever $\tilde{T}(s)>u_{\alpha}$. Then its FWER at $u \in \mathbb{R}$ is

$$
\operatorname{FWER}_{\tilde{T}}(u)=\mathbb{P}\left(\sup _{s \in \mathcal{H}_{0}} \tilde{T}(s)>u\right) .
$$

Defining $\tilde{T}_{0}=\sqrt{N}\left(\hat{\mu}_{N}(s)-\tilde{\mu}(s)\right) / \hat{\sigma}_{N}(s)$, then any $u_{\alpha}$ satisfying

$$
\mathbb{P}\left(\max _{s \in \mathcal{M}} \tilde{T}_{0}(s)>u_{\alpha}\right) \leq \alpha
$$

for $\alpha \in(0,1)$, controls the FWER in the strong sense at the level $\alpha$ because

$$
\operatorname{FWER}_{\tilde{T}}\left(u_{\alpha}\right) \leq \mathbb{P}\left(\max _{s \in \mathcal{M}_{\nu}} \tilde{T}_{0}(s)>u_{\alpha}\right) \leq \alpha
$$

for all $\mathcal{H}_{0} \subseteq \mathcal{M}_{\mathcal{V}}$ and $\tilde{T}_{0}(s)=\tilde{T}(s)$ for all $s \in \mathcal{H}_{0}$. To find such a threshold $u_{\alpha}$ we use the EC approximation (Adler [1981], Adler and Taylor [2007], Taylor et al. [2005]) to the excursion probability. In particular, letting $M_{u}\left(\tilde{T}_{0}\right)$ be the number of local maxima of $\tilde{T}_{0}$ over $\mathcal{M}_{\mathcal{V}}$ that lie above the level $u$ and $\chi_{\tilde{T}_{0}}(u)$ be the EC of the excursion set $\left\{s \in \mathcal{M}_{\mathcal{V}} \mid \tilde{T}_{0}(s)>u\right\}$ we have that

$$
\begin{equation*}
\mathbb{P}\left(\max _{s \in \mathcal{M}_{\nu}} \tilde{T}_{0}(s)>u\right) \leq \mathbb{E}\left[M_{u}\left(\tilde{T}_{0}\right)\right] \approx \mathbb{E}\left[\chi_{\tilde{T}_{0}}(u)\right]=\sum_{d=0}^{D} \mathcal{L}_{d} \rho_{d}^{\tilde{T}_{0}}(u) . \tag{23}
\end{equation*}
$$

Here $\rho_{d}^{\tilde{T}_{0}}$ are the EC densities of a (centered) $t$-field given, for example, in [Tay, 2007, p.915]. In order to control the FWER in the strong sense to a level $\alpha \in(0,1)$, we can thus find the largest $u_{\alpha}$ such that $\sum_{d=0}^{D} \hat{\mathcal{L}}_{d} \rho_{d}^{T}\left(u_{\alpha}\right)=\alpha$. At high thresholds $u_{\alpha}$ the number of local maxima is either zero or one and so $\mathbb{E}\left[\chi_{\tilde{T}_{0}}\left(u_{\alpha}\right)\right]$ is an extremely good approximation to $\mathbb{E}\left[M_{u_{\alpha}}\left(\tilde{T}_{0}\right)\right]$. Lower values of $\alpha$ yield higher thresholds $u_{\alpha}$, thus in practice where it is typical to take $\alpha \leq 0.05$, we expect the approximation in (23) to be accurate.

Traditional RFT inference in neuroimaging [Worsley et al., 1992, 1996, Tay, 2007] uses the same framework but only evaluates the fields on the lattice $\mathcal{V}$ and uses the LKC estimators given in Forman et al. [1995], Kie [1999] or Tay [2007] which are based on discrete derivatives. More precisely, for each $n \in\{1, \ldots, N\}$ it takes data $X_{n}$ on a lattice $\mathcal{V}$ (corresponding to the centers of voxels making up the brain), smoothes it with a kernel $K$ to obtain $\left\{\tilde{X}_{n}(v)\right.$ : $v \in \mathcal{V}\}$ and rejects all $v \in \mathcal{V}$ such that $\tilde{T}(v)>u_{\alpha}$ where $u_{\alpha}$ is obtained from the GKF approximation such that $\mathbb{P}\left(\max _{s \in \mathcal{M}} \tilde{T}_{0}(s)>u_{\alpha}\right) \approx \alpha$. Here $\mathcal{M} \subset \mathbb{R}^{3}$ represents for example

[^2]the brain, but has never been defined precisely in the literature as it probably was assumed to be unnecessary by the good lattice assumption. By construction this leads to valid, but conservative inference since $\max _{v \in \mathcal{V}} \tilde{T}_{0}(s) \leq \max _{s \in \mathcal{M}} \tilde{T}_{0}(s)$ and $u_{\alpha}$ approximates the tails of the distribution of $\max _{s \in \mathcal{M}} \tilde{T}_{0}(s)$ at level $\alpha$ and not the tails of the distribution of $\max _{v \in \mathcal{V}} \tilde{T}_{0}(v)$. SuRFs allow to remove this conservativeness by specifying $\mathcal{M}$ to be the voxel manifold $\mathcal{M}_{\mathcal{V}}$ and testing $\tilde{T}(s)>u_{\alpha}$ for all $s \in \mathcal{M}_{\mathcal{V}}$ meaning that $u_{\alpha}$ approximates the quantiles of the test statistic. Consequentially, our SuRF framework has strong FWER control over $\mathcal{M}_{\mathcal{V}}$ at level $\alpha$ up to the approximation in (23) and thus has a higher power than traditional RFT.

Effect Localization In our proposed framework the discrete data has been smoothed. Therefore precise localization of significant effects, i.e., finding $s \in S$ such that $\tilde{\mu}(s) \neq 0$, is only possible for the smoothed signal; still weaker localization results for the $v \in \mathcal{V}$ such that $\mu(v) \neq 0$ hold. ${ }^{4}$ The key observation is that for any $s \in \mathcal{M}_{\mathcal{V}}$ it holds that

$$
\begin{equation*}
\tilde{\mu}(s)>0 \Longleftrightarrow \exists v \in \mathcal{V} \cap \operatorname{supp}(K(s, \cdot)): \mu(v)>0 \tag{24}
\end{equation*}
$$

provided that $K(s, v) \geq 0$ for all $s \in \mathcal{M}_{\mathcal{V}}$ and all $v \in \mathcal{V}$. Thus, if we reject $\mathbf{H}_{0, s}$ we can conclude by (24) that there is at least one $v \in \mathcal{V} \cap \operatorname{supp}(K(s, \cdot))$ such that $\mu(v)>0$ and the strong control on the hypotheses (22) yields

$$
\begin{aligned}
\mathbb{P}\left(\left\{s^{\prime} \in S \mid \mathbf{H}_{0, s^{\prime}} \text { is rejected }\right\} \subseteq \mathcal{M}_{\mathcal{V}} \backslash \mathcal{H}_{0}\right) & =1-\mathbb{P}\left(\exists s \in \mathcal{H}_{0}: \mathbf{H}_{0, s} \text { is rejected }\right) \\
& \geq 1-\alpha .
\end{aligned}
$$

Consequentially, $\tilde{T}(s)>u_{\alpha}$ implies that the probability of incorrectly claiming that somewhere $\mu(v)>0$ on $\mathcal{V} \cap \operatorname{supp}(K(s, \cdot))$ is at most $\alpha$ which is a weaker form of FWER control on the signal $\mu$ than controlling it in the strong sense. However, it is stronger than controlling the FWER in the weak sense, as long as $\operatorname{supp}(K(s, \cdot)) \neq S$ for at least one $s \in S$, and as the signal-to-noise ratio typically is increased by applying smoothing, this type of FWER control on the hypotheses $\mathbf{H}_{0, v}: \mu(v) \leq 0$ vs. $\mathbf{H}_{1, v}: \mu(v)>0$ for $v \in \mathcal{V}$ satisfies the following "Heisenberg's uncertainty principle":

The smaller the supports of $K(\cdot, s), s \in S$, the better the test localizes effects, yet at the cost of losing power. On the other hand the higher the power to detect effects, i.e., the larger the support of $K(\cdot, s), s \in S$, the less precise the test localizes an effect.
Remark 9. Losing strong control with respect to $\mu$ due to smoothing the data is natural in applications such as fMRI. Smoothing is needed to increase the low signal-to-noise ratio and it is even debatable whether it is plausible to talk about the BOLD signal at a single voxel in a realistic fMRI experiment as the analyzed BOLD signal at $v$ is a distorted version of the observed data due to an extensive preprocessing pipeline which includes, among other aspects, motion correction and warping to a standardized brain.

## 4 Simulations

In this section we first use Monte-Carlo simulations to compare the performance of the SuRF estimator of the LKCs from Section 3.5 to existing LKC estimators; namely, the Hermite projection estimator (HP) and its bootstrap improvement (bHP) from [Telschow et al., 2023]

[^3]and the LKC estimators developed for stationary processes from Kie [1999] and Forman et al. [1995] which are used in established software, e.g., SPM and FSL. We abbreviate the latter two estimators as Kie and For respectively. We do not compare to the warping estimator Tay [2007] since the estimates, although computed differently, are almost identical to the estimates of the bHP, see [Telschow et al., 2023].

Our second set of Monte-Carlo simulations shows that our SuRF framework controls accurately the FWER in mass-univariate testing, while the traditional approach implemented in software such as SPM and FSL is conservative. In part 2 we even show that an adaptation of our methodology to non-Gaussian data is non-conservative on a gold-standard data set consisting of 7.000 resting state experiments from the UK Biobank.

### 4.1 Simulation Setup

We will generate our simulations by smoothing the Gaussian random field $\{X(v): v \in \mathcal{V}\}$ where $X(v)$ are i.i.d. $\mathcal{N}(0,1)$ distributed and $\mathcal{V} \subset \mathbb{R}^{D}$. To transform $X$ into a SuRF we smooth this discrete data by using the isotropic Gaussian kernel

$$
\begin{equation*}
K_{f}(x, v)=\left(\frac{4 \log (2)}{\pi f^{2}}\right)^{D / 2} e^{\frac{4 \log (2)\|x-v\|^{2}}{f^{2}}}, \quad f \in\{1,2,, \ldots, 6\}, \tag{25}
\end{equation*}
$$

which we parametrize by its full width at half maximum (FWHM) $f$ as is common in neuroimaging. We demonstrate the performance of different LKC estimators on a standardized almost stationary SuRF and a standardized non-stationary SuRF, which we define below.


Figure 2: Illustration of generating an almost stationary or non-stationary SuRF from a white noise field on a grid using $K_{3}$. The almost stationary field results from a convolution with the kernel $K_{3}$ before restricting the domain of the field. This circumvents boundary effects (top row). The non-stationary field (bottom row) results from first restricting to $\mathcal{V}_{2}$ and then convolving with the kernel $K_{3}$. The edge effects produces a non-stationary field.

The advantage of starting from the field $X$ which is i.i.d. on the lattice is that we can determine the theoretical LKCs for the normalized almost stationary and the normalized nonstationary SuRF quickly with high precision on a computer since the double sum in the Riemannian metric induced by the SuRF reduces to a single sum due to the lack of correlation, compare Proposition 7 and (7).

For the almost stationary simulations, for $D \in\{1,2,3\}, a \geq 0$, we take $\mathcal{V}=\mathcal{V}_{D}^{a} \subset \mathbb{R}^{D}$, where $\mathcal{V}_{1}^{a}=[1-a, 100+a] \cap \mathbb{Z}, \mathcal{V}_{2}^{a}=[1-a, 20+a]^{2} \cap \mathbb{Z}^{2}$ and $\mathcal{V}_{3}^{a}=[1-a, 20+a]^{3} \cap \mathbb{Z}^{3}$. The almost stationary SuRF is given by $\left(\tilde{X}, X, K_{f}, \mathcal{V}_{D}^{a}\right)$ with $a=\sqrt{2} f / \sqrt{\log (2)}$ and $\tilde{X}$ where
we choose $\mathcal{M}_{\mathcal{V}_{D}^{0}}$ to be the domain of $\tilde{X}$. We expand $\mathcal{V}_{D}^{0}$ by $a$ in each direction and restrict $\tilde{X}$ to the voxel manifold $\mathcal{M}_{\mathcal{V}_{D}^{0}}$ to remove boundary effects which allows a comparison with LKC estimators for stationary fields. We call this SuRF almost stationary because it is a stationary field on $\mathcal{V}_{D}^{0}$, but non-stationary on $\mathcal{M}_{\mathcal{V}_{D}^{0}}$. However, for $f$ larger than $\approx 2$ the LKCs of this field are almost identical to the LKCs of a stationary, unit variance Gaussian random field defined on $\mathcal{M}_{\mathcal{V}_{D}^{0}}$ with the covariance function $c\left(s, s^{\prime}\right)=K_{f}\left(s, s^{\prime}\right)$ with $s, s^{\prime} \in \mathcal{M}_{\mathcal{V}_{D}^{0}}$. The RFTtoolbox Davenport and Telschow [2023] can used to verify this as the theoretical LKCs of such fields can be approximated with high precision, see Appendix A.4. It is also supported by our observation in Davenport et al. that the SuRF LKC estimators yield precise FWHM estimates even at $f \approx 2.5$ where state-of-the-art methods are biased. The LKCs for stationary, unit-variance random fields are well known, see Worsley et al. [1996], and smoothing i.i.d. white noise with a Gaussian kernel is the typical way that FWER and LKC estimators have been validated [Hayasaka and Nichols, 2003, Taylor et al., 2007]. In practice, however, the data will be non-stationary and so this setting is unrealistic. To show that our methods are robust to non-stationarity we consider simulations where the fields are non-stationary. This is achieved by the initial lattice and deliberately not correcting for boundary effects. Precisely, we let $\mathcal{V}_{D}=\mathcal{V}_{D}^{0}$, for $D \in\{1,2,3\}$ where

$$
\begin{aligned}
& \mathcal{V}_{1}=([1,100] \cap \mathbb{Z}) \backslash\{2,4,8,9,11,15,20,21,22,40, \ldots, 45,60,62,64,65,98, \ldots, 100\} \\
& \mathcal{V}_{2}=\left\{x \in \mathbb{R}^{2} \mid x_{1} \in\{1,2,19,20\} \vee x_{2} \in\{1,2,19,20\}\right\} \cap[1,20]^{2} \\
& \mathcal{V}_{3}=\left\{x \in \mathbb{R}^{3} \mid x_{1} \in\{1,2,19,20\} \vee x_{2} \in\{1,2,19,20\} \vee x_{3} \in\{1,2,19,20\}\right\} \cap[1,20]^{3}
\end{aligned}
$$

and use $\mathcal{M}_{\mathcal{V}_{D}}$ as the domain of the SuRFs. This actually takes advantage of the boundary effect to produce non-stationary random fields. An illustration of individual sample fields in the case that $D=2$ is illustrated in Figure 2. In each of the simulation settings later, given a sample size $N \in \mathbb{N}$, we generate SuRFs $\tilde{X}_{1}, \ldots, \tilde{X}_{N}$ of SuRFs from $X_{1}, \ldots, X_{N} \sim X$.

### 4.2 Results of the LKC Estimation

For resolutions $r \in\{1,3,5\}$ and $D=2$ we compare the SuRF LKC estimator, see Section 3.5 , to the HP, bHP, Kiebel and Forman estimator obtained from samples of SuRFs evaluated on the grids with added resolution $r$. The results for $D=1,2$ are qualitatively similar, see Appendix A. In each setting, we run 1000 simulations in which we generate $N \in\{20,50,100\}$ SuRFs, compare Section 4.1, and estimate the LKCs from these fields.

Figure 3 shows boxplots of the resulting LKC estimates for $f=3$ and varying sample sizes $N \in\{20,50,100\}$ in the almost stationary setting. Figure 4 contains the same results for the non-stationary setting. At all resolutions the SuRF estimator seems to be unbiased and has a lower variance than the other estimators. Only the bHP estimator is comparable efficient, however, it has a small bias at resolution $r=1$. The HP estimator has a similar bias and a much larger variance, compare also Telschow et al. [2023]. The Kiebel and Forman estimators are biased for our non-stationary SuRF example and are even biased for the almost stationary SuRF for small values of $f$. This bias has been observed in the literature, see e.g., Kie [1999]. In Figure 5 and 6 we illustrate how the performance of the estimates of $\mathcal{L}_{2}$ depends on the smoothing bandwidth $f$. The corresponding plots for $\mathcal{L}_{1}$ can be found in Figures 12 and 13 in Appendix A. The SuRF estimator is the only approach which correctly estimates the LKCs for $f \leq 3 ; f=3$ is typical in neuroimaging. For the SuRF estimator only for $f=1$ a resolution increase of $r \geq 3$ is necessary to have unbiased estimates. Remarkably, for $f=2$ a resolution increase of $r=1$ seems to be sufficient for unbiased estimation using the SuRF


Figure 3: 2D Simulation results for estimation of the LKCs of the almost stationary SuRF described in Section 4.1. The smoothing bandwidth is $f=3$.


Figure 4: 2D Simulation results for estimation of the LKCs of the non-stationary SuRF described in Section 4.1. The smoothing bandwidth is $f=3$.
estimator. Furthermore, the SuRF estimator is 10 times faster to compute than its only reliable competitor the bHP estimator, see Table 1. Here we compare the average computation time of the SuRF, the bHPE and the Kiebel estimator on the stationary box example with $f=3$ and $N=100$ at different added resolution.

|  | SuRF |  |  | bHPE |  |  | Kiebel |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $r$ | 1 | 3 | 5 | 1 | 3 | 5 | 1 | 3 | 5 |
| $D=1$ | 0.06 | 0.05 | 0.05 | 7.01 | 13.35 | 16.13 | 4 | 0.03 | 0.04 |
| $D=2$ | 0.28 | 0.66 | 1.66 | 19.96 | 30.67 | 46.67 | 0.11 | 0.42 | 0.97 |
| $D=3$ | 18.22 | 80.45 | 239.30 | 196.44 | - | - | 8.24 | 47.90 | 151.75 |

Table 1: Computation time of LKC estimators in the stationary box example ( $N=100, f=3$ ). We show averages in seconds of 100 runs of the estimators. The times for resolution increases beyond $r=1$ and $D=3$ of the bHPE is not reported as they are very long.


Figure 5: 2D Simulation results for estimation of the LKCs of SuRFs derived from the stationary box example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.


Figure 6: 2D Simulation results for the estimation of the LKCs of SuRFs in the non-stationary setting. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.

### 4.3 Results of the FWER Simulation

Here we demonstrate that the SuRF framework improves upon traditional RFT [Worsley et al., 1996] in terms of FWER. To do so we calculate the FWER obtained from using RFT, see Section
3.6, from thresholding one-sample $t$-fields obtained from Gaussian random fields generated as described in Section 4.1. In each simulation setting, for $1 \leq b \leq B$ and $N \in\{20,50,100\}$ we obtain $t$-fields $T_{N, b, r}$ of resolution $r \in\{0,1, \infty\}$. Here $r=0$ corresponds to the traditional RFT approach, i.e., $T_{N, b, 0}$ is the test statistic described in Section 3.6 evaluated on the lattice $\mathcal{V}$ on which the original data is observed. Similarly, the case $r=1$ corresponds to $T_{N, b, 1}$ being the test statistic evaluated on $\mathcal{M}_{\mathcal{V}}^{(1)}$, i.e., $\mathcal{V} \cap \mathcal{M}_{\mathcal{V}}$ with an added resolution of 1 , compare (18). The case $r=\infty$ corresponds to the use of a SuRF, i.e., $T_{N, b, \infty}$ is the test statistic on $\mathcal{M}_{\mathcal{V}}$. In order to evaluate the FWER in the latter case we use numerical optimization methods (in particular sequential quadratic programming, Nocedal and Wright [2006]), initialized at the largest peaks of $T_{N, b, 1}$ to find the global maximum of $T_{N, b, \infty}$ over $\mathcal{M}_{\mathcal{V}}$. In each simulation we estimate the LKCs as described in Section 3.6 from using the sampling $\mathcal{M}_{\mathcal{V}}^{(1)}$ of $\mathcal{M}_{\mathcal{V}}$. Using $r=1$ for the estimation of the LKCs is justified by results of Section 4.2. In order to control the FWER at a level $\alpha=0.05$, we use the estimated LKCs to obtain an $\alpha$-level threshold $\hat{u}_{N, b}$ as described in Section 3.6. For $r \in\{0,1, \infty\}$ and $\mathcal{M}_{\mathcal{V}}^{(0)}=\mathcal{V}$ and $\mathcal{M}_{\mathcal{V}}^{(\infty)}=\mathcal{M}_{\mathcal{V}}$, the FWER is estimated by evaluating

$$
\frac{1}{B} \sum_{b=1}^{B} 1\left[\sup _{s \in \mathcal{M}_{v}^{(r)}} T_{N, b, r}(s)>\hat{u}_{N, b}\right],
$$

The results for $D=2,3$ are shown in Figures 7 and 8 respectively. In each simulation we also count the number of local maxima $l_{N, b, r}$ of $T_{N, b, r}$ on $\mathcal{M}_{\mathcal{V}}^{(r)}$ which exceed the level $\hat{u}_{N, b}$. This number approximates well the EC of the excursion set for $\alpha \approx 0.05$. Thus,

$$
\widehat{\mathrm{EEC}}=\frac{1}{B} \sum_{b=1}^{B} l_{N, b, r}
$$

plotted in yellow in the figures, is a good indicator of unbiasedness of our EEC estimate.
These results reproduce the well-known observation that the traditional $(r=0)$ approach is conservative, while our proposed the SuRF approach controls the FWER at the nominal rate in most settings. The difference is particularly observable at low smoothness levels, but even at $\mathrm{FWHM}=6$, evaluations of the random field on the original lattice are conservative. This effect is slightly more pronounced in 3D. Moving from the original $(r=0)$ lattice to the resolution 1 lattice already reduces the conservativeness which demonstrates that the main cause of the conservativeness is the mismatch between the discreteness of the data and modeling it as a random field over a manifold.

As can be seen from the yellow curve in the plots the EEC is typically well estimated. Only at the lowest considered smoothness level the EEC is underestimated, causing a slight conservativeness in the SuRF approach at $\mathrm{FWHM}=1$. This conservativeness could also be a consequence of the poor performance at $\mathrm{FWHM}=1$ of the SuRF LKC estimator with added resolution of 1, compare for example the first plot in Figures 5 and 6.


Figure 7: FWER results for the almost-stationary (top row) and non-stationary (bottom row) settings for $D=3$.


Figure 8: FWER results for the almost-stationary (top row) and non-stationary (bottom row) settings for $D=3$.

## 5 Discussion

This article solves the long-standing problem that voxelwise inference using RFT is conservative (Nichols and Hayasaka [2003], Taylor et al. [2007], Eklund et al. [2016]). Our solution relies on the distinction between data of an experiment and atoms of a probabilistic theory. This allowed us to identify the main causes of the conservativeness: (i) the test statistic is only evaluated on a discrete set given by the data although the atoms of RFT are random fields over a WS manifold and (ii) continuously available quantities are replaced by discrete counterparts. Therefore we used kernel smoothers as ferrymen transferring discrete data into random fields over a WS manifold - the atoms of RFT and showed in simulations that our SuRF methodology controls the FWER accurately at the given significance level. Although we stated our SuRF methodology for Gaussian related fields, this assumption can be relaxed. As demonstrated in Nardi et al. [2008], Telschow and Schwartzman [2022] the GKF approximates the EEC of the excursion sets of the test statistic $\tilde{T}$ well, if $\tilde{T}$ is asymptotically Gaussian. Hence for large enough sample sizes the SuRF methodology is applicable even under non-Gaussianity as the SuRF LKC estimates are consistent so long as derivatives of the sample covariance of the residuals converge uniformly to the derivatives of the covariance function of the limiting field, compare Telschow et al. [2023] for these type of arguments. However, the sample size required to deal with the non-Gaussianity in real fMRI experiments is relatively large as we demonstrate in part 2 . Here we also extend the SuRF methodology to non-Gaussian data and provide realistic simulations validating the non-conservativeness using a large resting state data set from the UK Biobank.

The importance of our findings for neuroimaging is twofold. Firstly, our GKF based method is computationally faster than resampling based inference such as permutation tests which control the FWER directly over the grid $\mathcal{V}$ and therefore yield FWER tests at level $\alpha$ at the cost of a high computational burden. Secondly, solving the conservativeness that has long caused power problems for voxelwise inference using RFT is a first step towards identifying and solving the problems of false positive rates in cluster inference (Eklund et al. [2016]) because the latter also relies on the GKF and applies continuous theory to smoothed fields evaluated on the voxel lattice [Chumbley and Friston, 2009].

As the estimator for $\mathcal{L}_{1}$ in 3D inspired by Theorem 2 is difficult to implement, we used a local stationary approximation which relies only on the first integral in (16). This does not influence voxelwise inference, even if the fields are highly non-stationary, because the estimation of $u_{\alpha}$ from the EEC is primarily driven by the values of $\hat{\mathcal{L}}_{2}$ and $\hat{\mathcal{L}}_{3}$. This can be seen using the typical value $u_{0.95} \approx 4.2$ from neuroimaging data (e.g., Supplementary material of Telschow et al. [2023]). It holds that $\rho_{3}^{\tilde{T}}\left(u_{\alpha}\right) / \rho_{1}^{\tilde{T}}\left(u_{\alpha}\right) \approx 2.6$ and $\rho_{2}^{\tilde{T}}\left(u_{\alpha}\right) / \rho_{1}^{\tilde{T}}\left(u_{\alpha}\right) \approx 1.6$ for sample size 50 and that typically $\mathcal{L}_{1} \ll \mathcal{L}_{3}$, compare our simulations and for real fMRI data the Supplementary material of Telschow et al. [2023]. The latter can be seen theoretically in the case of a stationary random field with square covariance function with bandwidth parameter $h$ over a convex WS manifold $\mathcal{M}$. Here $\mathcal{L}_{3}$ is the volume of $\mathcal{M}$ divided by $h^{3}, \mathcal{L}_{2}$ is half the surface area of $\mathcal{M}$ divided by $h^{2}$ and $\mathcal{L}_{1}$ is twice the diameter of $\mathcal{M}$ divided by $h$ [Worsley et al., 2004, Table 2]. Here the FWHM needs to be transformed into a bandwidth by multiplication with $\sqrt{4 \log (2)} / 2 \pi$ to transform resels (a concept from Worsley et al. [1996] designed for isotropic, stationary processes) into LKCs. Additionally, our simulations from this article and part 2 demonstrate that, the local stationary approximation of $\mathcal{L}_{1}$ does not influence the false positive rate. For cluster-size inference $\mathcal{L}_{1}$ may be more relevant as it has a greater effect at lower thresholds.

Last but not least, we want to emphasize that (i) and (ii) are relevant whenever continuous probabilistic theory is applied to discrete data. As such our proposed methodology based on

SuRFs is not restricted to voxelwise inference in brain imaging. In fact, the ideas presented apply equally when considering the coverage of confidence bands or the family-wise error rate of tests in functional data analysis where smoothing is applied for example via basis function representations or kernel smoothers.

There are a number of further scenarios in which the SuRF methodology is applicable. A concrete example is Davenport et al. [2022]. Here confidence regions for peaks of the signal of random fields over bounded open domains of $\mathbb{R}^{D}$ are developed and convolution fields are used to localize peaks of activation in fMRI and MEG. In fact, MEG is a natural domain in which our theory applies, because the power spectrum can be written as a convolution field, see the Supplementary material of Davenport et al. [2022] for details. A further potential application of the SuRF paradigm is to provide coverage probability excursion (CoPE) sets which attain the nominal coverage rate. CoPE sets provide confidence sets for the excursion above a value $c \in \mathbb{R}$ of a real-valued target function defined over a domain in $\mathbb{R}^{D}, D>0$, from noisy data, [Sommerfeld et al., 2018, Bowring et al., 2019, 2021]. These works found in simulations that the empirical coverage of CoPE sets is typically larger than the specified coverage, but converges to the correct coverage probability as the domain is sufficiently densely sampled. Theoretically, this has been explained in Theorem 1.b) from [Sommerfeld et al., 2018] and is illustrated in Figure 3 of Bowring et al. [2019]. Using the SuRFs in this setting can fix this overcoverage and yield CoPE sets that attain the nominal rate of coverage.

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## Appendix

## A Additional Simulation Results and Tables

## A. 1 LKC estimation for $D=1$



Figure 9: 1D Simulation results of estimation of the LKCs of the two SuRFs derived from the stationary box example and the non-stationary sphere example. The FWHM is $f=3$.



Figure 10: 1D Simulation results of estimation of the LKCs of SuRFs derived from the stationary box example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.


Figure 11: 1D Simulation results of estimation of the LKCs of SuRFs derived from the nonstationary sphere example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.

## A. 2 LKC estimation for $D=2$



Figure 12: 2D Simulation results for estimation of the LKCs of the almost stationary SuRF described in Section 4.1. The figures show the dependence of the LKC estimation on the smoothing bandwidth $f$ for sample size $N=100$.


Figure 13: 2D Simulation results of estimation of the LKCs of SuRFs derived from the nonstationary sphere example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.

## A. 3 LKC estimation for $D=3$



Figure 14: 3D Simulation results of estimation of the LKCs of the SuRFs derived from the stationary box example. The FWHM is $f=3$.


Figure 15: 3D Simulation results of estimation of the LKCs of the two SuRFs derived from the non-stationary sphere example. The FWHM is $f=3$. Note that the theoretical value for $\mathcal{L}_{1}$ is the theoretical value for the locally stationary $\mathcal{L}_{1}$. The true value for $\mathcal{L}_{1}$ is currently infeasible to obtain.


Figure 16: 3D Simulation results of estimation of the LKCs of SuRFs derived from the almost stationary box example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.


Figure 17: 3D Simulation results of estimation of the LKCs of SuRFs derived from the almost stationary box example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.


Figure 18: 3D Simulation results of estimation of the LKCs of SuRFs derived from the almoststationary sphere example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.


Figure 19: 3D Simulation results of estimation of the LKCs of SuRFs derived from the nonstationary sphere example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$. Note that the theoretical value for $\mathcal{L}_{1}$ is the theoretical value for the locally stationary $\mathcal{L}_{1}$. The true value for $\mathcal{L}_{1}$ is currently infeasible to obtain.


Figure 20: 3D Simulation results of estimation of the LKCs of SuRFs derived from the nonstationary sphere example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.


Figure 21: 3D Simulation results of estimation of the LKCs of SuRFs derived from the nonstationary sphere example. The results show the dependence of the LKC estimation on the FWHM used in the smoothing kernel for sample size $N=100$.

## A. 4 Comparison of the LKCs of the Almost Stationary Box Example and its Stationary Counterpart

Here we give an example of approximating the theoretical LKCs (up to an arbitrary resolution increase to approximate the integrals) of a SuRF derived from a random vector $(X(v): v \in V)$ satisfying $X\left(v_{1}\right)$ being independent of $X\left(v_{2}\right)$ for all $v_{1} \neq v_{2} \in \mathcal{V}$. This can be achieved using the function LKC_wncfield_theory() from the RFTtoolbox. We apply it to demonstrate that the theoretical LKCs from the almost stationary box example used in our simulations which has covariance function $\tilde{\mathfrak{c}}_{f}$ depending on $f>0$ has essentially the same LKCs for $f$ larger than $\approx 2.5$ as the zero-mean Gaussian field with covariance function $\mathfrak{c}_{f}(x, y) \exp \left(-4 \log (2)|x-y|^{2} / f^{2}\right)$. The LKCs of the latter random field over a domain $S \subset \mathbb{R}^{D}, D \in\{1,2,3\}$, are given by

$$
\begin{array}{ll}
D=1: & \mathcal{L}_{1}=\frac{\operatorname{vol}(S)}{\sqrt{4 \log (2)} f} \\
D=2: & \mathcal{L}_{1}=\frac{\operatorname{vol}(\partial S)}{2 \sqrt{4 \log (2)} f}, \quad \mathcal{L}_{2}=\frac{\operatorname{vol}(S)}{4 \log (2) f^{2}} \\
D=3: & \mathcal{L}_{1}=\frac{2 \operatorname{Diameter}(S)}{\sqrt{4 \log (2)} f}, \quad \mathcal{L}_{2}=\frac{\operatorname{vol}(\partial S)}{2 \cdot 4 \log (2) f^{2}}, \quad \mathcal{L}_{3}=\frac{\operatorname{vol}(S)}{f^{3}(4 \log (2))^{3 / 2}},
\end{array}
$$

compare for example Telschow et al. [2023] and in particular Table 2 from Worsley et al. [2004] for $D=3$. Note that the given closed form of $\mathcal{L}_{1}$ requires that $S$ is also convex, Worsley et al. [2004]. To approximate the theoretical LKCs we use an added resolution of 11 for $D \in\{1,2\}$ and for the sake of computation time use an added resolution of 7 for $D=3$. The results are presented in Tables 2-4.

| $D=1$ | $f$ | 1 | 1.5 | 2 | 2.5 | 3 | 3.5 | 4 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathcal{L}_{1}$ | $\tilde{\mathfrak{c}}_{f}$ | 146.52 | 110.41 | 83.25 | 66.60 | 55.50 | 47.57 | 41.63 |
|  | $\mathfrak{c}_{f}$ | 166.51 | 111.01 | 83.26 | 66.60 | 55.50 | 47.57 | 41.63 |

Table 2: Comparison of the theoretical LKCs for different smoothing bandwidths between the almost stationary box example for $D=1$, which has the covariance function $\tilde{\mathfrak{c}}_{f}$, and the zeromean stationary Gaussian field having covariance $\tilde{\mathfrak{c}}_{f}$ over the domains specified in Section 4.1 of the main manuscript.

| $D=2$ | $f$ | 1 | 1.5 | 2 | 2.5 | 3 | 3.5 | 4 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathcal{L}_{1}$ | $\tilde{\mathfrak{c}}_{f}$ | 58.61 | 44.16 | 33.30 | 26.64 | 22.20 | 19.03 | 16.65 |
|  | $\mathfrak{c}_{f}$ | 66.60 | 44.40 | 33.30 | 26.64 | 22.20 | 19.03 | 16.65 |
| $\mathcal{L}_{2}$ | $\tilde{\mathfrak{c}}_{f}$ | 858.72 | 487.59 | 277.24 | 177.45 | 123.23 | 90.53 | 69.31 |
|  | $\mathfrak{c}_{f}$ | 1109.00 | 492.90 | 277.26 | 177.45 | 123.23 | 90.53 | 69.31 |

Table 3: Comparison of the theoretical LKCs for different smoothing bandwidths between the almost stationary box example for $D=2$, which has the covariance function $\tilde{\mathfrak{c}}_{f}$, and the zeromean stationary Gaussian field having covariance $\tilde{\mathfrak{c}}_{f}$ over the domains specified in Section 4.1 of the main manuscript.

| $D=3$ | $f$ | 1 | 1.5 | 2 | 2.5 | 3 | 3.5 | 4 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathcal{L}_{1}$ | $\tilde{\mathfrak{c}}_{f}$ | 87.91 | 66.24 | 49.95 | 39.96 | 33.30 | 28.54 | 24.98 |
|  | $\mathfrak{c}_{f}$ | 99.91 | 66.60 | 49.95 | 39.96 | 33.30 | 28.54 | 24.98 |
| $\mathcal{L}_{2}$ | $\tilde{\mathfrak{c}}_{f}$ | 2576.13 | 1462.77 | 831.72 | 532.34 | 369.68 | 271.60 | 207.94 |
|  | $\mathfrak{c}_{f}$ | 3327.11 | 1478.71 | 831.78 | 532.34 | 369.68 | 271.6 | 207.94 |
| $\mathcal{L}_{3}$ | $\tilde{\mathfrak{c}}_{f}$ | 25163.37 | 10766.66 | 4616.20 | 2363.73 | 1367.90 | 861.42 | 577.08 |
|  | $\mathfrak{c}_{f}$ | 36933.30 | 10943.20 | 4616.66 | 2363.73 | 1367.90 | 861.42 | 577.08 |

Table 4: Comparison of the theoretical LKCs for different smoothing bandwidths between the almost stationary box example for $D=3$, which has the covariance function $\tilde{\mathfrak{c}}_{f}$, and the zeromean stationary Gaussian field having covariance $\tilde{\mathfrak{c}}_{f}$ over the domains specified in Section 4.1 of the main manuscript.

## B LKCs induced by a normalized field

## B. 1 Induced Riemannian metric of a normalized random field

The most important quantity for the Gaussian Kinematic formula is the Riemannian metric induced by a random field. It is the backbone of the GKF for Gaussian related fields developed in Taylor et al. [2006].

In this section $f$ denotes a zero-mean random field with almost surely continuously, differentiable sample paths over the domain $\overline{\mathcal{M}}$ and we call the random field $f / \sqrt{\operatorname{Var}[f]}$ the normalized field derived from $f$. Recall that a vector field $V \in \mathcal{T} \overline{\mathcal{M}}$ can be interpreted as a first order differential operator, i.e., for all $h \in C^{1}(\overline{\mathcal{M}})$ the expression $V h: \overline{\mathcal{M}} \rightarrow \mathbb{R}, s \mapsto(V h)(s)$ defines a function in $C(\overline{\mathcal{M}})$. This can be made precise by taking a local chart $\varphi: \overline{\mathcal{M}} \supseteq \mathcal{U} \rightarrow \varphi(\mathcal{U})=\mathcal{W} \subset \mathbb{R}^{D}$ of $\overline{\mathcal{M}}$ with inverse $\psi=\varphi^{-1}$. The vector field $V$ in local coordinates $(\mathcal{U}, \varphi)$ can be written as

$$
\begin{equation*}
V=\sum_{d=1}^{D} V_{d} \partial_{d}, \quad V_{1}, \ldots, V_{D} \in C(\mathcal{W}) \tag{26}
\end{equation*}
$$

where $\partial_{d}$ is the vector field on $\mathcal{W}$ defined by

$$
\begin{equation*}
\partial_{d}(h \circ \psi)(x)=\frac{\partial h \circ \psi}{\partial x_{d}}(x)=\partial_{d}^{x} h(\psi(x)), \quad d=1, \ldots, D \tag{27}
\end{equation*}
$$

where $h \in C(\mathcal{U})$ and $x \in \varphi(\mathcal{U})$.
Definition 5. Let $V, W \in \mathcal{T} \overline{\mathcal{M}}$ be differentiable vector fields and $f$ unit-variance random field over a manifold $\overline{\mathcal{M}}$ with almost surely differentiable sample paths. Then

$$
\begin{equation*}
\bar{\Lambda}_{s}(V, W)=\operatorname{Cov}[V f(s), W f(s)], s \in \overline{\mathcal{M}} \tag{28}
\end{equation*}
$$

is called the induced Riemannian metric of $f$ on $\overline{\mathcal{M}}$. In local coordinates at a point $z \in \mathcal{W}$ it is represented by

$$
\begin{align*}
\bar{\Lambda}_{z}(V, W) & =\left.\sum_{d, d^{\prime}=1}^{D} V_{d} W_{d^{\prime}} \partial_{d}^{x} \partial_{d^{\prime}}^{y} \operatorname{Cov}[f(\psi(x)), f(\psi(y))]\right|_{(x, y)=(z, z)}  \tag{29}\\
& =\sum_{d, d^{\prime}=1}^{D} V_{d} W_{d^{\prime}} \operatorname{Cov}\left[\left.\partial_{d}^{x} f(\psi(x))\right|_{x=z},\left.\partial_{d^{\prime}}^{y} f(\psi(y))\right|_{y=z}\right] .
\end{align*}
$$

Here $V_{1}, \ldots, V_{D}, W_{1}, \ldots, W_{D} \in C(\mathcal{W})$ are the coordinate coefficients representing the vector fields $V$ and $W$.

Remark 10. Assumption (G2) from the main manuscript ensures that this is a Riemannian metric on $\overline{\mathcal{M}}$.

Since the vector fields $\partial_{d}, d=1, \ldots, D$, form a basis of $\mathcal{T U}$ the Riemannian metric induced by the random field $f$ can be written in local coordinates as the $D \times D$ matrix having components

$$
\begin{equation*}
\bar{\Lambda}_{d d^{\prime}}(z)=\operatorname{Cov}\left[\left.\partial_{d}^{x} f(\psi(x))\right|_{x=z},\left.\partial_{d^{\prime}}^{y} f(\psi(y))\right|_{y=z}\right] \tag{30}
\end{equation*}
$$

For simplicity in what follows, we establish the following alternative notations suppressing the dependencies on $f$ :

$$
\begin{equation*}
\operatorname{Cov}\left[\partial_{d}^{x} f(x), f(y)\right]=\left\langle\partial_{d}^{x}, 1_{y}\right\rangle, \quad \operatorname{Cov}\left[\partial_{d}^{x} f(x), \partial_{d^{\prime}}^{y} f(y)\right]=\left\langle\partial_{d}^{x}, \partial_{d^{\prime}}^{y}\right\rangle \tag{31}
\end{equation*}
$$

Similarly, $\left\|1_{x}\right\|^{2}=\operatorname{Var}[f(x)]$ and $\left\|\partial_{d}^{x}\right\|^{2}=\operatorname{Var}\left[\partial_{d}^{x} f(x)\right]$.

Theorem 5. The Riemannian metric on $\overline{\mathcal{M}}$ induced by a normalized random field is given in local coordinates by

$$
\begin{equation*}
\bar{\Lambda}_{d d^{\prime}}(x)=\frac{\left\langle\partial_{d}^{x}, \partial_{d^{\prime}}^{x}\right\rangle}{\left\|1_{x}\right\|}-\frac{\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle}{\left\|1_{x}\right\|^{2}} \tag{32}
\end{equation*}
$$

The computation of Lipshitz-Killing curvatures requires the shape operator and the Riemannian curvature. Therefore, the next theorem computes the Christoffel symbols of the first kind, which can be used to express these quantities in local coordinates. This is because most fundamental geometric quantities such as the covariant derivative and the Riemannian curvature are functions of the Christoffel symbols and their derivatives.

Theorem 6. The Christoffel symbols $\bar{\Gamma}_{k d d^{\prime}}, k, d, d^{\prime} \in\{1, \ldots, D\}$ of the first kind of the induced Riemannian metric by a normalized field are given by

$$
\begin{align*}
\bar{\Gamma}_{k d d^{\prime}}(x)= & \frac{\left\langle\partial_{k}^{x} \partial_{d}^{x}, \partial_{d^{\prime}}^{x}\right\rangle}{\left\|1_{x}\right\|}-\frac{\left\langle\partial_{k}^{x} \partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle}{\left\|1_{x}\right\|^{2}}-\frac{\left\langle\partial_{k}^{x}, \partial_{d^{\prime}}^{x}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle}{\left\|1_{x}\right\|^{2}} \\
& -\frac{\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, \partial_{d^{\prime}}^{x}\right\rangle}{\left\|1_{x}\right\|^{2}}+2 \frac{\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, f_{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle}{\left\|1_{x}\right\|^{3}} \tag{33}
\end{align*}
$$

In terms of Christoffel symbols the covariant derivative $\bar{\nabla}$ on $\overline{\mathcal{M}}$ is expressed in the local chart $\varphi$ by

$$
\bar{\nabla}_{\partial_{d}} \partial_{d^{\prime}}=\sum_{h=1}^{D} \sum_{h^{\prime}=1}^{D} \bar{\Lambda}^{h h^{\prime}} \bar{\Gamma}_{d d^{\prime} h^{\prime}} \partial_{h}=\bar{\Lambda}^{-1}\left(\bar{\Gamma}_{d d^{\prime} 1}, \ldots, \bar{\Gamma}_{d d^{\prime} D}\right)^{T}
$$

Here $\bar{\Lambda}^{d d^{\prime}}$ denotes the $\left(d, d^{\prime}\right)$-entry of the inverse of $\bar{\Lambda}$ in the coordinates $(\mathcal{U}, \varphi)$. This formula can be used to extend the covariant derivative to any vector field, i.e.,

$$
\begin{equation*}
\bar{\nabla}_{V} W=\sum_{d=1}^{D} \sum_{d^{\prime}=1}^{D} V_{d} \bar{\nabla}_{\partial_{d}} W_{d^{\prime}} \partial_{d^{\prime}}=\sum_{d=1}^{D} \sum_{d^{\prime}=1}^{D} V_{d}\left(\partial_{d}\left(W_{d^{\prime}}\right) \partial_{d^{\prime}}+W_{d^{\prime}} \bar{\nabla}_{\partial_{d}} \partial_{d^{\prime}}\right) \tag{34}
\end{equation*}
$$

In particular, if all $\Gamma_{d d^{\prime} d^{\prime \prime}}=0, d, d^{\prime}, d^{\prime \prime} \in\{1, \ldots, D\}$, then

$$
\begin{equation*}
\bar{\nabla}_{V} W=\sum_{d=1}^{D} \sum_{d^{\prime}=1}^{D} V_{d} \partial_{d}\left(W_{d^{\prime}}\right) \partial_{d^{\prime}} \tag{35}
\end{equation*}
$$

This happens, if $\bar{\Lambda}(s)=\bar{\Lambda}\left(s^{\prime}\right)$ for all $s, s^{\prime} \in \overline{\mathcal{M}}$.
The last geometric quantity required to compute LKCs is the Riemannian curvature tensor $\bar{R}$. The curvature tensor in local coordinates can be written as the tensor having entries

$$
\begin{equation*}
\bar{R}_{l k d d^{\prime}}^{\varphi}=\partial_{l} \bar{\Gamma}_{k d d^{\prime}}-\partial_{k} \Gamma_{l d d^{\prime}}+\sum_{m, n=1}^{D}\left(\bar{\Gamma}_{l d m} \bar{\Lambda}^{m n} \bar{\Gamma}_{k d^{\prime} n}-\bar{\Gamma}_{k d m} \bar{\Lambda}^{m n} \bar{\Gamma}_{l d^{\prime} n}\right) \tag{36}
\end{equation*}
$$

for $i, j, k, l \in\{1, \ldots, D\}$, compare [Adler and Taylor, 2007, eq. (7.7.4)]. By now, we derived almost all quantities to state the Riemannian curvature tensor in local coordinates. The missing quantities are the derivatives of the Christoffel symbols which can be found in the next Lemma.

Lemma 1. The difference of the derivatives of Christoffel symbols in the Riemannian curvature tensor (36) of a normalized field can be expressed as:

$$
\begin{aligned}
& \partial_{l} \bar{\Gamma}_{k d d^{\prime}}(x)-\partial_{k} \bar{\Gamma}_{l d d^{\prime}}(x) \\
& =\frac{\left\langle\partial_{k}^{x} \partial_{d}^{x}, \partial_{l}^{x} \partial_{d^{\prime}}^{x}\right\rangle-\left\langle\partial_{l}^{x} \partial_{d}^{x}, \partial_{k}^{x} \partial_{d^{\prime}}^{x}\right\rangle}{\left\|1_{x}\right\|} \\
& -\left\|1_{x}\right\|^{-2}\left(\left\langle\partial_{l}^{x}, 1_{x}\right\rangle\left\langle\partial_{k}^{x} \partial_{d}^{x}, \partial_{d^{\prime}}^{x}\right\rangle-\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{l}^{x} \partial_{d}^{x}, \partial_{d^{\prime}}^{x}\right\rangle\right. \\
& +\left\langle\partial_{k}^{x} \partial_{d}^{x}, \partial_{l}^{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle-\left\langle\partial_{l}^{x} \partial_{d}^{x}, \partial_{k}^{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle \\
& +\left\langle\partial_{k}^{x} \partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{l}^{x} \partial_{d^{\prime}}^{x}, 1_{x}\right\rangle-\left\langle\partial_{l}^{x} \partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{k}^{x} \partial_{d^{\prime}}^{x}, 1_{x}\right\rangle \\
& +\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{l}^{x} \partial_{d^{\prime}}^{x}, \partial_{d}^{x}\right\rangle-\left\langle\partial_{l}^{x}, 1_{x}\right\rangle\left\langle\partial_{k}^{x} \partial_{d^{\prime}}^{x}, \partial_{d}^{x}\right\rangle \\
& +\left\langle\partial_{l}^{x}, \partial_{d}^{x}\right\rangle\left\langle\partial_{k}^{x}, \partial_{d^{\prime}}^{x}\right\rangle-\left\langle\partial_{k}^{x}, \partial_{d}^{x}\right\rangle\left\langle\partial_{l}^{x}, \partial_{d^{\prime}}^{x}\right\rangle \\
& +\left\langle\partial_{l}^{x}, \partial_{d}^{x}\right\rangle\left\langle\partial_{k}^{x}, \partial_{d^{\prime}}^{x}\right\rangle-\left\langle\partial_{k}^{x}, \partial_{d}^{x}\right\rangle\left\langle\partial_{l}^{x}, \partial_{d^{\prime}}^{x}\right\rangle \\
& \left.+\left\langle\partial_{k}^{x}, \partial_{l}^{x} \partial_{d^{\prime}}^{x}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle-\left\langle\partial_{l}^{x}, \partial_{k}^{x} \partial_{d^{\prime}}^{x}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\right) \\
& +2\left\|1_{x}\right\|^{-3}\left(\left\langle\partial_{l}^{x}, 1_{x}\right\rangle\left\langle\partial_{k}^{x} \partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle-\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{l}^{x} \partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle\right. \\
& +\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, \partial_{l}^{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle-\left\langle\partial_{l}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, \partial_{k}^{x}\right\rangle\left\langle\partial_{d^{\prime}}^{x}, 1_{x}\right\rangle \\
& +\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{l}^{x} \partial_{d^{\prime}}^{x}, 1_{x}\right\rangle-\left\langle\partial_{l}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x},\right\rangle\left\langle\partial_{k}^{x} \partial_{d^{\prime}}^{x}, 1_{x}\right\rangle \\
& \left.+\left\langle\partial_{l}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{k}^{x}, \partial_{d^{\prime}}^{x}\right\rangle-\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\left\langle\partial_{l}^{x}, \partial_{d^{\prime}}^{x}\right\rangle\right)
\end{aligned}
$$

Proof. Simple, but lengthy computation.
To compute the first LKC $\mathcal{L}_{1}$ for a 3-dimensional manifold the trace of the Riemannian tensor is needed which in the coordinates $(\mathcal{U}, \varphi)$ can be expressed in terms of the entries of the inverse of the square root of the Riemannian metric and the Riemannian tensor as

$$
\begin{align*}
\operatorname{tr}(\bar{R})=\sum_{i, j, k, l=1}^{3} \bar{R}_{i j k l}^{\varphi}( & \frac{\bar{\Lambda}_{1 i}^{-1 / 2} \bar{\Lambda}_{1 j}^{-1 / 2} \bar{\Lambda}_{1 k}^{-1 / 2} \bar{\Lambda}_{1 l}^{-1 / 2}}{2} \\
& +\frac{\bar{\Lambda}_{2 i}^{-1 / 2} \bar{\Lambda}_{2 j}^{-1 / 2} \bar{\Lambda}_{2 k}^{-1 / 2} \bar{\Lambda}_{2 l}^{-1 / 2}}{2} \\
& +\frac{\bar{\Lambda}_{3 i}^{-1 / 2} \bar{\Lambda}_{3 j}^{-1 / 2} \bar{\Lambda}_{3 k}^{-1 / 2} \bar{\Lambda}_{3 l}^{-1 / 2}}{2}  \tag{37}\\
& +\bar{\Lambda}_{1 i}^{-1 / 2} \bar{\Lambda}_{2 j}^{-1 / 2} \bar{\Lambda}_{1 k}^{-1 / 2} \bar{\Lambda}_{2 l}^{-1 / 2} \\
& +\bar{\Lambda}_{1 i}^{-1 / 2} \bar{\Lambda}_{3 j}^{-1 / 2} \bar{\Lambda}_{1 k}^{-1 / 2} \bar{\Lambda}_{3 l}^{-1 / 2} \\
& \left.+\bar{\Lambda}_{2 i}^{-1 / 2} \bar{\Lambda}_{3 j}^{-1 / 2} \bar{\Lambda}_{2 k}^{-1 / 2} \bar{\Lambda}_{3 l}^{-1 / 2}\right)
\end{align*}
$$

## B. 2 Definition of LKCs

Lipshits Killing Curvatures $\mathcal{L}_{1}, \ldots, \mathcal{L}_{D}$ are the intrinsic volumes of a compact $D$-dimensional Whitney stratified manifold $(\mathcal{M}, \Lambda)$ isometrically embedded into $(\overline{\mathcal{M}}, \bar{\Lambda})$. Here $\Lambda$ and $\bar{\Lambda}$ denote Riemannian metrics of $\mathcal{M}$ and $\overline{\mathcal{M}}$. They are related by $\bar{\Lambda} \circ \iota=\Lambda$ where $\iota: \mathcal{M} \rightarrow \overline{\mathcal{M}}$ is the embedding and hence $\mathcal{M}$ is isometrically embedded into $\overline{\mathcal{M}}$. In this section we make the
formula for the LKCs [Adler and Taylor, 2007, Definition 10.7.2] for the LKCs of an at most 3-dimensional Whitney stratified manifold more explicit. From the definition of the LKCs it is easy to deduce

$$
\begin{equation*}
\mathcal{L}_{D-1}=\operatorname{vol}_{D-1}\left(\partial_{D-1} \mathcal{M}\right) \quad \text { and } \quad \mathcal{L}_{D}=\operatorname{vol}_{D}\left(\partial_{D} \mathcal{M}\right) . \tag{38}
\end{equation*}
$$

Here $\partial_{d} \mathcal{M}$ denotes the $d$-dimensional stratum of $\mathcal{M}$ and the volume is the volume of the Riemannian manifold $\left(\partial_{d} \mathcal{M},\left.\Lambda\right|_{\partial_{d} \mathcal{M}}\right)$ where $\left.\Lambda\right|_{\partial_{d} \mathcal{M}}$ is the restriction of $\Lambda$ to $\partial_{d} \mathcal{M}$.

For a 3 -dimensional manifold it remains to compute $\mathcal{L}_{1}$, which we derive from [Adler and Taylor, 2007, Definition 10.7.2] and some geometric computations in the next proposition.

Proposition 8. Let $(\overline{\mathcal{M}}, \bar{\Lambda})$ be a closed Riemannian manifold of dimension 3 and $\mathcal{M} \subset \overline{\mathcal{M}}$ be a compact Whitney stratified manifold of dimension 3 isometrically embedded into $\overline{\mathcal{M}}$. Then

$$
\begin{align*}
\mathcal{L}_{1}=\frac{1}{2 \pi} & \int_{\partial_{1} \mathcal{M}} \int_{\mathbb{S}\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \alpha(v) \mathcal{H}_{1}(d v) \mathcal{H}_{1}(d s) \\
& +\frac{1}{2 \pi} \int_{\partial_{2} \mathcal{M}} \bar{\Lambda}\left(\bar{\nabla}_{U_{s}} U_{s}+\bar{\nabla}_{V_{s}} V_{s}, N_{s}\right) \mathcal{H}_{2}(d s)  \tag{39}\\
& -\frac{1}{2 \pi} \int_{\partial_{3} \mathcal{M}} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}(\bar{R}) \mathcal{H}_{3}(d s)
\end{align*}
$$

Here $\mathcal{H}_{1}(d v)$ is the volume form induced on the sphere

$$
\mathbb{S}\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)=\left\{v \in \mathcal{T}_{s} \overline{\mathcal{M}} \mid \bar{\Lambda}(v, v)=1 \wedge \bar{\Lambda}(v, w)=0 \text { for } w \in \mathcal{T}_{s} \partial_{1} \mathcal{M}\right\}
$$

by $\bar{\Lambda}$ and $\mathcal{H}_{d}(d s)$ the volume form of $\partial_{d} \mathcal{M}$. Moreover, $\alpha(v)$ denotes the normal Morse index given in [Adler and Taylor, 2007, Scn 9.2.1] and $U, V, N$ is a piecewise differentiable vector field on $\partial_{2} \mathcal{M}$ such that $U_{s}, V_{s}$ form an orthonormal basis for $\mathcal{T}_{s} \partial_{2} \mathcal{M}$ for all $s \in \partial_{2} \mathcal{M}$ and $N$ is outward pointing normal vector field.

Remark 11. In the special case that the metric $\bar{\Lambda}$ is constant, it holds that

$$
\begin{aligned}
\mathcal{L}_{1}= & \frac{1}{2 \pi} \int_{\partial_{1} \mathcal{M}} \int_{\mathbb{S}\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \alpha(v) \mathcal{H}_{1}(d v) \mathcal{H}_{1}(d s) \\
& +\frac{1}{2 \pi} \int_{\partial_{2} \mathcal{M}} \bar{\Lambda}\left(\bar{\nabla}_{U_{s}} U_{s}+\bar{\nabla}_{V_{s}} V_{s}, N_{s}\right) \mathcal{H}_{2}(d s),
\end{aligned}
$$

since the curvature tensor $\bar{R}$ vanishes. If $\mathcal{M}=\mathcal{M}_{\mathcal{V}}$ is a voxel manifold than even

$$
\mathcal{L}_{1}=\frac{1}{2 \pi} \int_{\partial_{1} \mathcal{M}} \int_{\mathbb{S}\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \alpha(v) \mathcal{H}_{1}(d v) \mathcal{H}_{1}(d s),
$$

as $\nabla_{U_{s}} U_{s}=\nabla_{V_{s}} V_{s}=0$ on $\partial_{2} \mathcal{M}$ by (35) and $\bar{\Lambda}$ being constant.

## C Proofs of Results in the Appendix

## C. 1 Proof of Theorem 5

Proof. Interchanging expectation and derivatives yields

$$
\begin{align*}
\bar{\Lambda}_{d d^{\prime}}(z)= & \left.\partial_{d}^{x} \partial_{d^{\prime}}^{y} \frac{\operatorname{Cov}\left[f_{x}, f_{y}\right]}{\left\|1_{x}\right\|\left\|1_{y}\right\|}\right|_{(x, y)=(z, z)} \\
= & \frac{\left\langle\partial_{d}^{x}, \partial_{d^{\prime}}^{y}\right\rangle}{\left\|1_{x}\right\|\left\|1_{y}\right\|}-\frac{\left\langle\partial_{d}^{x}, 1_{y}\right\rangle\left\langle 1_{y}, \partial_{d^{\prime}}^{y}\right\rangle}{\left\|1_{x}\right\|\left\|1_{y}\right\|^{3}}  \tag{40}\\
& -\frac{\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\left\langle 1_{x}, \partial_{\left.d^{d^{\prime}}\right\rangle}^{\left\|1_{x}\right\|^{3}\left\|1_{y}\right\|}+\left.\frac{\left\langle 1_{x}, 1_{y}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\left\langle 1_{y}, \partial_{d^{\prime}}^{y}\right.}{\left\|1_{x}\right\|^{3}\left\|1_{y}\right\|^{3}}\right|_{(x, y)=(z, z)}\right.}{} \\
= & \frac{\left\langle\partial_{d}^{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{2}}-\frac{\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}}
\end{align*}
$$

## C. 2 Proof of Theorem 6

Proof. From the definition of the Christoffel symbols one can derive that

$$
\begin{equation*}
\Gamma_{k d d^{\prime}}(z)=\left.\frac{\partial^{3} \tilde{\mathfrak{c}}(x, y)}{\partial_{x_{k}} \partial_{x_{d}} \partial_{y_{d^{\prime}}}}\right|_{(x, y)=(z, z)}, \tag{41}
\end{equation*}
$$

see also [Adler and Taylor, 2007, eq. 12.2.17]. Here $\tilde{\mathfrak{c}}(x, y)=\frac{\operatorname{Cov}\left[f_{x}, f_{y}\right]}{\sqrt{\operatorname{Cov}\left[f_{x}, f_{x}\right]} \sqrt{\operatorname{Cov}\left[f_{y}, f_{y}\right]}}$. Thus, by simply taking another derivate of (40), we obtain

$$
\begin{aligned}
\Gamma_{k d d^{\prime}}(z)= & \partial_{k}^{x} \partial_{d}^{x} \partial_{d^{\prime}}^{y} \frac{\operatorname{Cov}\left[f_{x}, f_{y}\right]}{\left.\sqrt{\left\|1_{x}\right\|\left\|1_{y}\right\|}\right|_{(x, y)=(z, z)}} \\
= & \frac{\left\langle\partial_{k}^{x} \partial_{d}^{x}, \partial_{d^{\prime}}^{x}\right\rangle}{\left\|1_{x}\right\|^{2}}-\frac{\left\langle\partial_{k}^{x}, 1_{x}\right\rangle\left\langle\partial_{d}^{x}, \partial_{d^{\prime}}^{x}\right\rangle}{\left\|1_{x}\right\|^{4}}-\partial_{k}^{x} \frac{\left\langle\partial_{d}^{x}, 1_{y}\right\rangle\left\langle 1_{y}, \partial_{d^{\prime}}^{y}\right\rangle}{\left\|1_{x}\right\|\left\|1_{y}\right\|^{3}} \\
& -\partial_{k}^{x} \frac{\left.x \partial_{d}^{x}, 1_{x}\right\rangle\left\langle 1_{x}, \partial_{d^{\prime}}^{y}\right\rangle}{\left\|1_{x}\right\|^{3}\left\|1_{y}\right\|}+\left.\partial_{k}^{x} \frac{\left\langle 1_{x}, 1_{y}\right\rangle\left\langle\partial_{d}^{x}, 1_{x}\right\rangle\left\langle 1_{y}, \partial_{d^{\prime}}^{y}\right\rangle}{\left\|1_{x}\right\|^{3}\left\|1_{y}\right\|^{3}}\right|_{(x, y)=(z, z)} \\
= & \frac{\left\langle\partial_{k}^{z} \partial_{d}^{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{2}}-\frac{\left\langle\partial_{k}^{z}, 1_{z}\right\rangle\left\langle\partial_{d}^{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}}-\frac{\left\langle\partial_{k}^{z} \partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}} \\
& +\frac{\left\langle\partial_{k}^{z}, 1_{z}\right\rangle\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{6}}-\frac{\left\langle\partial_{k}^{z} \partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle+\left\langle\partial_{d}^{z}, \partial_{k}^{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}} \\
& -\frac{\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle\partial_{k}^{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}}+3 \frac{\left\langle\partial_{k}^{z}, 1_{z}\right\rangle\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{6}} \\
& +\frac{\left\langle\partial_{k}^{z}, 1_{z}\right\rangle\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{6}}+\frac{\left\langle\partial_{k}^{z} \partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}}+\frac{\left\langle\partial_{d}^{z}, \partial_{k}^{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}} \\
& -3 \frac{\left\langle\partial_{k}^{z}, 1_{z}\right\rangle\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{6}} \\
= & \frac{\left\langle\partial_{k}^{z} \partial_{d}^{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{2}}-\frac{\left\langle\partial_{k}^{z}, 1_{z}\right\rangle\left\langle\partial_{d}^{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}}-\frac{\left\langle\partial_{k}^{z} \partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}} \\
& -\frac{\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle\partial_{k}^{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{4}}+2 \frac{\left\langle\partial_{k}^{z}, 1_{z}\right\rangle\left\langle\partial_{d}^{z}, 1_{z}\right\rangle\left\langle 1_{z}, \partial_{d^{\prime}}^{z}\right\rangle}{\left\|1_{z}\right\|^{6}}
\end{aligned}
$$

## C. 3 Proof of Proposition 8

Proof. Using $D-d^{\prime}=D^{\prime}$ we have that the LKCs of a Whitney stratified manifold $\mathcal{M}$ are defined by

$$
\begin{align*}
\mathcal{L}_{d}= & \sum_{d^{\prime}=d}^{D} \frac{1}{(2 \pi)^{\frac{d^{\prime}-d}{2}}} \sum_{l=0}^{\left\lfloor\frac{d^{\prime}-d}{2}\right\rfloor} \frac{(-1)^{l} C\left(D^{\prime}, d-d^{\prime}-2 l\right)}{l!\left(d-d^{\prime}-2 l\right)!}  \tag{42}\\
& \times \int_{\partial_{d^{\prime}} \mathcal{M}} \int_{\mathbb{S}\left(\mathcal{T}_{s} \partial_{d^{\prime}} \mathcal{M}^{\perp}\right)} \operatorname{Tr}_{s} \mathcal{T}_{d^{\prime}} \mathcal{M}\left(R^{l} S_{\nu_{D^{\prime}}}^{d^{\prime}-d-2 l}\right) \alpha\left(\nu_{D^{\prime}}\right) \mathcal{H}_{-1}\left(d \nu_{D^{\prime}}\right) \mathcal{H}_{d^{\prime}}(d s)
\end{align*}
$$

compare [Adler and Taylor, 2007, Definition 10.7.2]. This formula requires further explanations. The constant $C(m, i)$ is defined in [Adler and Taylor, 2007, eq. (10.5.1), p.233], i.e.,

$$
C(m, i)=\left\{\begin{array}{ll}
\frac{(2 \pi)^{\frac{i}{2}}}{s_{m+i}}, & m+i>0,  \tag{43}\\
1, & m=0
\end{array} \quad \text { with } \quad s_{m}=\frac{2 \pi^{m / 2}}{\Gamma(m / 2)}\right.
$$

which implies $C(m, 0)=\Gamma(m / 2) / 2 / \pi^{m / 2}$. Moreover, $\mathcal{H}_{D-d-1}$ is the volume form on $\mathbb{S}\left(\mathcal{T}_{s} \partial_{d} \mathcal{M}^{\perp}\right)$ and $\mathcal{H}_{d}$ the volume form of $\partial_{d} \mathcal{M} . R$ denotes the Riemannian curvature tensor of the different
strata on $\mathcal{M}$ depending on the strata the integral integrates over. In particular, note that $\bar{R}=R$ for $\partial_{3} \mathcal{M}$.

From this $\mathcal{L}_{1}$ can be simplified as follows:

$$
\begin{aligned}
& \mathcal{L}_{1}=C(3,0) \int_{\partial_{1} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{1} \mathcal{M}}\left(R^{0} S_{\nu_{2}}^{0}\right) \alpha\left(\nu_{2}\right) \mathcal{H}_{1}\left(d \nu_{2}\right) \mathcal{H}_{1}(d s) \\
& +\frac{C(1,1)}{\sqrt{2 \pi}} \int_{\partial_{2} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{2} \mathcal{M}^{\perp}\right)} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{2} \mathcal{M}}\left(R^{0} S_{\nu_{1}}^{1}\right) \alpha\left(\nu_{1}\right) \mathcal{H}_{0}\left(d \nu_{1}\right) \mathcal{H}_{2}(d s) \\
& +\frac{C(0,2)}{4 \pi} \int_{\partial_{3} \mathcal{M}} \int_{S(\mathbb{Q})} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}\left(R^{0} S_{\nu_{0}}^{2}\right) \alpha\left(\nu_{0}\right) \mathcal{H}_{-1}\left(d \nu_{0}\right) \mathcal{H}_{3}(d s) \\
& -\frac{C(0,0)}{2 \pi} \int_{\partial_{3} \mathcal{M}} \int_{S(\mathbb{O})} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}\left(R^{1} S_{\nu_{0}}^{0}\right) \alpha\left(\nu_{0}\right) \mathcal{H}_{-1}\left(d \nu_{0}\right) \mathcal{H}_{3}(d s) \\
& =\frac{1}{2 \pi} \int_{\partial_{1} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \alpha\left(\nu_{2}\right) \mathcal{H}_{1}\left(d \nu_{2}\right) \mathcal{H}_{1}(d s) \\
& +\frac{1}{2 \pi} \int_{\partial_{2} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{2} \mathcal{M}^{\perp}\right)} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{2} \mathcal{M}}\left(1 \cdot S_{\nu_{1}}^{1}\right) \alpha\left(\nu_{1}\right) \mathcal{H}_{0}\left(d \nu_{1}\right) \mathcal{H}_{2}(d s) \\
& +\frac{1}{4 \pi} \int_{\partial_{3} \mathcal{M}} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}\left(1 \cdot S_{\mathscr{\oplus}}^{2}\right) \mathcal{H}_{3}(d s) \\
& -\frac{1}{2 \pi} \int_{\partial_{3} \mathcal{M}} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}\left(R^{1} S_{\oplus}^{0}\right) \mathcal{H}_{3}(d s) \\
& =\frac{1}{2 \pi} \int_{\partial_{1} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \alpha\left(\nu_{2}\right) \mathcal{H}_{1}\left(d \nu_{2}\right) \mathcal{H}_{1}(d s) \\
& +\frac{1}{2 \pi} \int_{\partial_{2} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{2} \mathcal{M}^{\perp}\right)} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{2} \mathcal{M}}\left(S_{\nu_{1}}^{1}\right) \alpha\left(\nu_{1}\right) \mathcal{H}_{0}\left(d \nu_{1}\right) \mathcal{H}_{2}(d s) \\
& -\frac{1}{2 \pi} \int_{\partial_{3} \mathcal{M}} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}(\bar{R}) \mathcal{H}_{3}(d s) \\
& =\frac{1}{2 \pi} \int_{\partial_{1} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \alpha\left(\nu_{2}\right) \mathcal{H}_{1}\left(d \nu_{2}\right) \mathcal{H}_{1}(d s) \\
& +\frac{1}{2 \pi} \int_{\partial_{2} \mathcal{M}} g\left(\bar{\nabla}_{e_{1}(s)} e_{1}(s), \tilde{\nu}(s)\right)+g\left(\bar{\nabla}_{e_{2}(s)} e_{2}(s), \tilde{\nu}(s)\right) \mathcal{H}_{2}(d s) \\
& -\frac{1}{2 \pi} \int_{\partial_{3} \mathcal{M}} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}(\bar{R}) \mathcal{H}_{3}(d s) \\
& =\frac{1}{2 \pi} \int_{\partial_{1} \mathcal{M}} \int_{S\left(\mathcal{T}_{s} \partial_{1} \mathcal{M}^{\perp}\right)} \alpha\left(\nu_{2}\right) \mathcal{H}_{1}\left(d \nu_{2}\right) \mathcal{H}_{1}(d s) \\
& +\frac{1}{2 \pi} \int_{\partial_{2} \mathcal{M}} g\left(\bar{\nabla}_{e_{1}(s)} e_{1}(s)+\bar{\nabla}_{e_{2}(s)} e_{2}(s), \tilde{\nu}(s)\right) \mathcal{H}_{2}(d s) \\
& -\frac{1}{2 \pi} \int_{\partial_{3} \mathcal{M}} \operatorname{Tr}^{\mathcal{T}_{s} \partial_{3} \mathcal{M}}(\bar{R}) \mathcal{H}_{3}(d s)
\end{aligned}
$$

Here $\tilde{\nu}(s)$ is the inward pointing normal at $x$ in $\partial_{2} \mathcal{M}$ and $e_{1}(s), e_{2}(s)$ an orthonormal basis of $\mathcal{T}_{s} \partial_{2} \mathcal{M}$ and used Remark (10.5.2) [Adler and Taylor, 2007, p.233], i.e.,

$$
S_{\mathbb{O}}^{j}= \begin{cases}1, & j=0  \tag{44}\\ 0, & \text { otherwise }\end{cases}
$$

## D Proofs of the Results in the Main Manuscript

## D. 1 Proof of Proposition 2

Proof. Define $q=p /(p-1)$ if $p>1$ and $q=\infty$ if $p=1$. Using the triangle inequality and Hölder's inequality yields for the charts $\left(\bar{U}_{\alpha}, \bar{\varphi}_{\alpha}\right), \alpha \in\{1, \ldots, P\}$, in the atlas of $\overline{\mathcal{M}}$ covering $\mathcal{M}$, all $x, y \in \bar{\varphi}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ that

$$
\begin{aligned}
\left|\tilde{X}_{\alpha}(x)-\tilde{X}_{\alpha}(y)\right| & =\left|\sum_{v \in \mathcal{V}}\left(K\left(\bar{\varphi}_{\alpha}^{-1}(x), v\right)-K\left(\bar{\varphi}_{\alpha}^{-1}(y), v\right)\right) X(v)\right| \\
& \leq \sqrt[q]{\sum_{v \in \mathcal{V}}\left|K\left(\bar{\varphi}_{\alpha}^{-1}(x), v\right)-K\left(\bar{\varphi}_{\alpha}^{-1}(y), v\right)\right|^{q}} \sqrt[p]{\sum_{v \in \mathcal{V}} X(v)^{p}} \\
& \leq \sqrt[q]{\sum_{v \in \mathcal{V}} A^{q}} \sqrt[p]{\sum_{v \in \mathcal{V}} X(v)^{p}}\|x-y\|^{\gamma} \\
& \leq L\|x-y\|^{\gamma} .
\end{aligned}
$$

Here $A$ bounds the Hölder constants of $K\left(\bar{\varphi}_{\alpha}^{-1}(\cdot), v\right)$, for all $\alpha \in\{1, \ldots, P\}$ and all $v \in \mathcal{V}$ from above, and $L=|\mathcal{V}| A \sqrt[p]{\sum_{v \in \mathcal{V}} X(v)^{p}}$. If $p=1$ then the statement with the $q$-th root is the maximum over $\mathcal{V}$ instead of the $q$-norm. The result follows as by assumption $\mathbb{E}\left[L^{p}\right]$ is finite.

## D. 2 Proof of Proposition 3

Proof. The functions $K\left(\bar{\varphi}_{\alpha}^{-1}(\cdot), v\right)$ are Lipshitz continuous for each $v \in \mathcal{V}$ since they are $\mathcal{C}^{1}$ and $\overline{\mathcal{M}}$ is compact. Because $\mathcal{V}$ is finite and $\alpha \in\{1, \ldots, P\}$, there exists an $M>0$ that bounds all the Lipschitz constants of the functions $K\left(\bar{\varphi}_{\alpha}^{-1}(\cdot), v\right)$. Thus, applying Proposition 2 with $p=2$ and $\gamma=1$ yields for the charts $\left(\bar{U}_{\alpha}, \bar{\varphi}_{\alpha}\right), \alpha \in\{1, \ldots, P\}$, in the atlas of $\overline{\mathcal{M}}$ covering $\mathcal{M}$, all $x, y \in \bar{\varphi}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ such that $0<\|x-y\|<1$ that

$$
\mathbb{E}\left[\left(\tilde{X}_{\alpha}(x)-\tilde{X}_{\alpha}(y)\right)^{2}\right] \leq \mathbb{E}\left[L^{2}\|x-y\|^{2}\right]=\mathbb{E}\left[L^{2}\right]\|x-y\|^{2} .
$$

The claim follows since $x^{2} \leq(\log |x|)^{-2}$ for $0<x<1$.

## D. 3 Proof of Proposition 4

Proof. As the property is local, we can w.l.o.g. assume that $\mathcal{M}$ is a compact domain in $\mathbb{R}^{D}$ and drop the chart notation for simplicity. Given $x \in \mathcal{M}$, suppose that there exist sets of real constants $a, a_{i}, a_{j k}, c(1 \leq i \leq D, 1 \leq j \leq k \leq D)$ such that

$$
a \tilde{X}(s)+\sum_{i=1}^{D} a_{i} \tilde{X}_{i}(s)+\sum_{1 \leq j \leq k \leq D} a_{j k} \tilde{X}_{j k}(s)=c,
$$

which implies that

$$
\begin{aligned}
a \sum_{v \in \mathcal{V}_{s}} K(s, v) X(v) & +\sum_{i=1}^{D} a_{i} \sum_{v \in \mathcal{V}_{s}} \partial_{i}^{s} K(s, v) X(v) \\
& +\sum_{1 \leq j \leq k \leq D} a_{j k} \sum_{v \in \mathcal{V}_{s}} \partial_{j k}^{s} K(s, v) X(v)=c .
\end{aligned}
$$

Non-degeneracy of $\left(X(v): v \in \mathcal{V}_{s}\right)$ then implies that for all $v \in \mathcal{V}_{s}$

$$
a K(s, v)+\sum_{i=1}^{D} a_{i} \partial_{i}^{s} K(s, v)+\sum_{1 \leq j \leq k \leq D} a_{j k} \partial_{j k}^{s} K(s, v)=0,
$$

which by the linear independence constraint implies that the constants are all zero. This proves non-degeneracy of $\left(Y(s), \nabla Y(s),\left(\nabla^{2} Y(s)\right)\right)$.

For the normalized field $\tilde{X} / \sqrt{\operatorname{Var}[\tilde{X}]}=\tilde{X} / \sigma$, we note that

$$
\nabla \frac{X}{\sigma}=\frac{\nabla X}{\sigma}-\frac{X \nabla \sigma}{\sigma^{2}}=\frac{\nabla X}{\sigma}-\frac{\nabla \sigma}{\sigma}\left(\frac{X}{\sigma}\right)
$$

and

$$
\nabla^{2} \frac{X}{\sigma}=\frac{\nabla^{2} X}{\sigma}-\frac{2(\nabla X)^{T}(\nabla \sigma)}{\sigma^{2}}-\frac{X \nabla^{2} \sigma}{\sigma^{2}}+\frac{2(\nabla \sigma)^{T}(\nabla \sigma) X}{\sigma^{3}} .
$$

Hence using an invertible matrix we can transform $\left(X(s), \nabla X(s), \mathbb{V}\left(\nabla^{2} X(s)\right)\right)$ into $\left(Z(s), \nabla Z(s), \mathbb{V}\left(\nabla^{2} Z(s)\right)\right)$. Thus, $\left(Z(s), \nabla Z(s), \mathbb{V}\left(\nabla^{2} Z(s)\right)\right)$ is non-degenerate by Lemma A. 2 from Davenport and Telschow [2022].

## D. 4 Proof of Proposition 6

## D.4.1 Establishing non-degeneracy of the isotropic kernel and its derivatives under linear transformations

Lemma 2. Suppose that $\mathcal{V}$ satisfies the conditions of Proposition 6 and let $K^{*}=e^{-\|s-v\|^{T} / 2}$ be the $D$-dimensional isotropic Gaussian kernel. Then given constants $c, a_{i}, a_{j k}$ (for $1 \leq i \leq D$ and $1 \leq j \leq k \leq D), s \in \mathbb{R}^{D}$ and an invertible symmetric matrix $\Omega^{\prime} \in \mathbb{R}^{D \times D}$ such that

$$
\begin{align*}
c K^{*}\left(\Omega^{\prime} s, \Omega^{\prime} v\right) & +\sum_{j=1}^{D} a_{j} \partial_{j}^{x} K^{*}\left(\Omega^{\prime} s, \Omega^{\prime} v\right)  \tag{45}\\
& +\sum_{1 \leq j \leq k \leq D} a_{j k} \partial_{j k}^{x} K^{*}\left(\Omega^{\prime} s, \Omega^{\prime} v\right)=0
\end{align*}
$$

for all $v \in \mathcal{V}$ then $c=a_{i}=a_{j k}=0$ for $1 \leq i \leq D$ and $1 \leq j \leq k \leq D$.
Proof. For all $v \in \mathcal{V}$, letting $s^{*}=\Omega^{\prime} s$ and dividing (45) by $e^{-\|s-v\|^{2} / 2}$, it follows that

$$
\begin{aligned}
c & +\sum_{j=1}^{D} a_{j}\left(s_{j}^{*}-\sum_{l=1}^{D} \Omega_{j l}^{\prime} v_{l}\right) \\
& +\sum_{1 \leq j \leq k \leq D} a_{j k}\left(\left(s_{j}^{*}-\sum_{l=1}^{D} \Omega_{j l}^{\prime} v_{l}\right)\left(s_{k}^{*}-\sum_{l=1}^{D} \Omega_{k l}^{\prime} v_{l}\right)-\delta_{j k}\right)=0 .
\end{aligned}
$$

In particular,

$$
\begin{align*}
c & +\sum_{j=1}^{D} a_{j}\left(s_{j}^{*}-\sum_{l=1}^{D} \Omega_{j l}^{\prime} v_{l}\right) \\
& +\sum_{1 \leq j, k \leq D} a_{j k}^{\prime}\left(\left(s_{j}^{*}-\sum_{l=1}^{D} \Omega_{j l}^{\prime} v_{l}\right)\left(s_{k}^{*}-\sum_{l=1}^{D} \Omega_{k l}^{\prime} v_{l}\right)-\delta_{j k}\right)=0 . \tag{46}
\end{align*}
$$

where $a_{j k}^{\prime}=a_{j k} / 2, j \neq k$ and $a_{j j}^{\prime}=a_{j j}$.
For $i \in\{1, \ldots, D\}$, fixing $\left(v_{1}, \ldots, v_{i-1}, v_{i+1}, \ldots, v_{D}\right)$, we can view (46) as a quadratic in $v_{i}$. As such the only way that it can have more than two distinct solutions is if where the coefficient of $v_{i}^{2}$ is zero, i.e.,

$$
\begin{equation*}
\sum_{1 \leq j, k \leq D} a_{j k}^{\prime} \Omega_{j i}^{\prime} \Omega_{k i}^{\prime}=\left(\Omega^{\prime} A^{\prime} \Omega^{\prime}\right)_{i i}=0 . \tag{47}
\end{equation*}
$$

Similarly the coefficient of $v_{i}$ must be zero, i.e.,

$$
\begin{equation*}
\sum_{j=1}^{D} a_{j} \Omega_{j i}^{\prime}+\sum_{j, k} \Omega_{j i}^{\prime}\left(s_{k}^{*}-\sum_{m \neq i} \Omega_{k m}^{\prime} v_{m}\right)+\sum_{j, k}\left(s_{k}^{*}-\sum_{l \neq i} \Omega_{j l}^{\prime} v_{l}\right) \Omega_{k i}^{\prime}=0 . \tag{48}
\end{equation*}
$$

Now allowing $v_{n}$ to vary for some $n \neq i$, by the same logic, the coefficient of $v_{n}$ in (48) is equal to zero, i.e.,

$$
\sum_{j, k} a_{j k}^{\prime} \Omega_{j i}^{\prime} \Omega_{k n}^{\prime}+\sum_{j, k} a_{j k}^{\prime} \Omega_{j n}^{\prime} \Omega_{k i}^{\prime}=\left(\Omega^{\prime} A^{\prime} \Omega^{\prime}\right)_{i n}+\left(\Omega^{\prime} A^{\prime} \Omega^{\prime}\right)_{n i}=2\left(\Omega^{\prime} A^{\prime} \Omega^{\prime}\right)_{i n}=0 .
$$

As such $\left(\Omega^{\prime} A^{\prime} \Omega^{\prime}\right)_{i n}=\left(\Omega^{\prime} A^{\prime} \Omega^{\prime}\right)_{n i}=0$ for all $i \neq n$. Combining this with (47), it follows that $\Omega^{\prime} A^{\prime} \Omega^{\prime}=0$. In particular $A^{\prime}=0$ as $\Omega^{\prime}$ is invertible. Thus, the remaining linear equation in $v_{i}$ from (46) can only have more than one solution if

$$
\sum_{j} a_{j} \Omega_{j i}^{\prime}=0
$$

Therefore $\left(\Omega^{\prime} a\right)_{i}=0$, where $a=\left(a_{1}, \ldots, a_{D}\right)^{T}$. Since this holds for all $i$ and $\Omega^{\prime}$ is invertible, we obtain $a=0$. Finally this implies that $c=0$.

## D.4.2 Establishing Proposition 6

Proof. We can write $K(s, v)=K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right)$ where $K^{*}=e^{-\|s-v\|^{2} / 2}$ is the isotropic Gaussian kernel. Arguing as in proof of Lemma 1 of Davenport and Telschow [2022] (taking $\phi(s)=\Sigma^{-1 / 2} s$ and $\varphi$ to be the identity in their notation), for each $s \in S$ we have

$$
\mathbb{V}\left(\nabla^{2} K(s, v)\right)=L\left(\Sigma^{1 / 2} \otimes \Sigma^{1 / 2}\right) R \mathbb{V}\left(\nabla^{2} K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right)\right)
$$

where $\nabla^{2}$ as usual always denotes the Hessian with respect to the first argument and $L \in$ $\mathbb{R}^{D(D+1) / 2 \times D^{2}}$ is the elimination matrix and $R \in \mathbb{R}^{D^{2} \times D(D+1) / 2}$ is the duplication matrix, the precise definitions of which can be found in Magnus and Neudecker [1980]. The matrix $L\left(\Sigma^{1 / 2} \otimes \Sigma^{1 / 2}\right) R$ is invertible by Lemma 4.4.iv of Magnus and Neudecker [1980], and the fact that $\Sigma^{1 / 2}$ is invertible.

Moreover $\nabla K\left(\Sigma^{-1 / 2} s, v\right)=\Sigma^{-1 / 2} \nabla K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right)$. As such there is an invertible linear transformation between the vector

$$
\left(K(s, v), \nabla K(s, v), \mathbb{V}\left(\nabla^{2} K(s, v)\right)\right)
$$

and the vector

$$
\left(K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right), \nabla K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right), \mathbb{V}\left(\nabla^{2} K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right)\right)\right) .
$$

In particular if there exists constants $c, a_{d}, a_{k l}$ such that $1 \leq d \leq D$ and $1 \leq k \leq l \leq D$ (with at least one of them being non-zero) such that

$$
\begin{equation*}
c K(s, v)+\sum_{d=1}^{D} a_{d} \partial_{d}^{s} K(s, v)+\sum_{1 \leq k \leq l \leq D} a_{k l} \partial_{k l}^{s} K(s, v)=0 \tag{49}
\end{equation*}
$$

then there existing corresponding constants $\left(c^{*}, a_{d}^{*}, a_{k l}^{*}\right)$ (with at least one of them being nonzero) such that

$$
\begin{align*}
c^{*} K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right) & +\sum_{d=1}^{D} a_{d}^{*} \partial_{d}^{s} K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right)  \tag{50}\\
& +\sum_{1 \leq k \leq l \leq D} a_{k l}^{*} \partial_{k l}^{s} K^{*}\left(\Sigma^{-1 / 2} s, \Sigma^{-1 / 2} v\right)=0 .
\end{align*}
$$

Applying Lemma 2 yields a contradiction and thus establishes the result.

## D. 5 Proof of Proposition 7

Proof. Note that for a $\operatorname{SuRF}(\tilde{X}, X, K, \mathcal{V})$ we obtain the following identity

$$
\begin{equation*}
\left\langle\partial_{d} K_{s}, K_{s^{\prime}}\right\rangle=\operatorname{Cov}\left[\partial_{d}^{s} \tilde{X}(s), \tilde{X}\left(s^{\prime}\right)\right]=\left\langle\partial_{d}^{s}, 1_{s^{\prime}}\right\rangle \tag{51}
\end{equation*}
$$

Hence the result for a normalized SuRF is a Corollary of Theorem 5.

## D. 6 Proof of Theorem 1

Proof. In order to apply Theorem 12.4.2 from Adler and Taylor [2007] we need to prove that the assumptions (G1)-(G3) hold. The assumption that $\tilde{X}_{\alpha}$ is Gaussian with almost surely $C^{2}$-sample paths clearly holds by the assumption that $X$ is a Gaussian field on $\mathcal{V}$ and $K(\cdot, v) \in$ $C^{3}(\overline{\mathcal{M}})$ for all $v \in \mathcal{V}$. The non-degeneracy condition follows from Proposition 4. The last assumption that there is an $\epsilon>0$ such that

$$
\mathbb{E}\left[\left(\partial_{d d^{\prime}} \tilde{X}_{\alpha}(x)-\partial_{d d^{\prime}} \tilde{X}_{\alpha}(y)\right)^{2}\right] \leq K|\log \|x-y\||^{-(1+\gamma)}
$$

for some $K>0$, all $d, d^{\prime} \in\{1, \ldots, D\}$ and for the charts $\left(\bar{U}_{\alpha}, \bar{\varphi}_{\alpha}\right), \alpha \in\{1, \ldots, P\}$, in the atlas of $\overline{\mathcal{M}}$ covering $\mathcal{M}$, all $x, y \in \bar{\varphi}\left(\bar{U}_{\alpha}\right) \cap \mathcal{M}$ such that $|x-y|<\epsilon$ is established in Proposition 3.

## D. 7 Proof of Theorem 2

Computation of $\Theta(x)=\int_{\mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M}_{\nu}^{\perp}\right)} \alpha(\nu) \mathcal{H}_{1}(d \nu)$ for Voxel Manifolds. In order to compute $\Theta(x)$ we need to introduce the normal Morse index $\alpha(\nu)$. We specialize here to the case of $\mathcal{M}_{\mathcal{V}}$ being a voxel manifold embedded into $\mathbb{R}^{3}$, yet the exact same concept is defined for any Whitney stratified manifold, compare [Adler and Taylor, 2007, Scn 9.2.1].

For any $x \in \mathcal{M}_{\mathcal{V}}$ and any direction $\nu \in \mathbb{S}\left(\mathcal{T}_{x} \mathcal{M}_{\mathcal{V}}\right)$ the normal Morse index is one minus the local Euler characteristic of the intersection of $\mathcal{M}_{\mathcal{V}}$, the $\delta$-ball centered at $x$ and the affine plane $\left\{\lambda \in \mathbb{R}^{3}: \lambda^{T} \nu+x+\epsilon=0\right\}$ for $\epsilon>0$. If $\delta$ is sufficiently small, then this Euler characteristic does not depend on $\epsilon$ provided that $\epsilon$ is small enough. Although this definition sounds complicated at first, it can be easily computed for all $x$ and $\nu$ for a voxel manifold. Note that $\mathcal{M}_{\mathcal{V}}$ or more
precisely an open neighbourhood of it might be endowed with a different Riemannian metric than the standard Riemannian metric on $\mathbb{R}^{3}$, which in our case is the induced Riemannian metric $\boldsymbol{\Lambda}$ by a unit-variance random field $f$. In this case $\nu \in \mathbb{S}\left(\mathcal{T}_{x} \mathbb{R}^{3}\right)$ are vectors $\nu \in \mathbb{R}^{3}$ such that $\nu^{T} \boldsymbol{\Lambda}(x) \nu=1$.

In the case that $x \in \partial_{3} \mathcal{M}_{\mathcal{V}}$ and $\nu \in \mathbb{S}\left(\mathcal{T}_{x} \mathbb{R}^{3}\right)$ it is obvious that $\alpha(\nu)=0$ because the intersection of the affine plain defined by $\nu$ and the $\delta$-ball is always homeomorphic to a filled disk, which has Euler characteristic 1. Similarly, if $x \in \partial_{2} \mathcal{M}_{\mathcal{V}}$, then $\alpha(\nu)=0$ for all $\nu \in$ $\mathbb{S}\left(\mathcal{T}_{x} \mathbb{R}^{3}\right) \backslash\left\{\nu_{\text {out }}\right\}$. Here $\nu_{\text {out }}$ is the unique outside pointing normal (w.r.t. the Euclidean metric!) at $x \in \partial_{2} \mathcal{M}_{\mathcal{V}}$ and it holds that $\alpha\left(\nu_{\text {out }}\right)=1$ as the plain defined by $\nu_{\text {out }}$ for a small enough $\delta$-ball is parallel to $\partial_{2} \mathcal{M}_{\mathcal{V}}$ and therefore its intersection with $\mathcal{M}_{\mathcal{V}}$ is empty. The important cases for us and the interesting cases happen at edges of the voxel manifold, i.e., for $x \in \partial_{1} \mathcal{M}_{\mathcal{V}}$. Here the behavior of $\alpha(\nu)$ can be classified by the three types of possible edges: the convex, the double convex and the concave edge. These cases are shown in Figure 1 in the main manuscript.

The behavior of $\alpha$ for directions $\nu \in \mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{1}{\mathcal{V}}\right)$ is demonstrated in Figure 22 within the hyperplane $x+\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{\perp}{\mathcal{V}}$ with $\mathcal{M}_{\mathcal{V}}$. Here we show the two possible intersection scenarios of the hyperplane (bold green line) orthogonal to the direction $-\nu \in \mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{1}{\mathcal{V}}\right)$ (green arrow) with a small $\delta$-ball (dotted black line) and the voxel manifold $\mathcal{M}_{\mathcal{\nu}}$ from the definition of $\alpha$. In particular, it can be seen that $\alpha(\nu)$ is constant, if $-\nu$ is inside the normal cone $\mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}$ (for a definition see [Adler and Taylor, 2007, p.189]) and constant on $\mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{1}{\mathcal{V}}\right) \backslash \mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}$ indpendent on the type of edge to which $x$ belongs. The geometric embedding of the the intersection of the (geometric) normal cone $x+\mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}$ with the hyperplane $x+\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{\downarrow}{\mathcal{V}}$ is represented by the red shaded areas. From this we deduce that on a convex edge the Euler characteristic of the intersection of the green hyperplane with the $\delta$-ball and $\mathcal{M} \mathcal{V}$ is 0 , if $-\nu \in \mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{\perp}{\mathcal{V}}\right)$, as the intersection is empty and 1 otherwise because the intersection is homeomorphic to a disk. Similar, it holds that the Euler characteristic of the intersection is 2 for $x$ on a double convex or concave edge, if $-\nu \in \mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{1}{\mathcal{V}}\right)$, as the intersection is homeomorphic to the disjoint union of two disks and 1 else because the intersection is homeomorphic to a disk. Therefore we obtain that $\alpha(\nu)$ for $\nu \in \mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{1}{\mathcal{V}}\right)$ is given by

$$
\alpha(\nu)=\left\{\begin{array}{rl}
1, & \text { if } x \text { on convex edge and }-\nu \in\left(\mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}\right)^{\circ}  \tag{52}\\
-1, & \text { if } x \text { on a double convex edge and }-\nu \in\left(\mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}\right)^{\circ} \\
-1, \text { if } x \text { on concave edge and }-\nu \in\left(\mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}\right)^{\circ} \\
0, \text { else }
\end{array} .\right.
$$

Using this we can compute the function $\Theta(x)$.
Lemma 3. Let $\mathcal{M}_{\mathcal{V}}$ be a voxel manifold, $x \in \partial_{1} \mathcal{M}_{\mathcal{V}}$ and $U_{x}, V_{x}, W_{x}$ denote an $O N$ frame for $\mathcal{T}_{x} \overline{\mathcal{M}_{\mathcal{V}}}$. Define $M=V_{x} \times W_{x}=\left(m_{1}(x), m_{2}(x), m_{3}(x)\right)^{T}$. Define

$$
\beta(x)=\arccos \left(\frac{m_{2}(x) m_{3}(x)}{\sqrt{m_{2}^{2}(x)+m_{1}^{2}(x)} \sqrt{m_{3}^{2}(x)+m_{1}^{2}(x)}}\right) .
$$

Then we obtain

$$
\Theta(x)= \begin{cases}\pi-\beta(x), & \text { if } x \text { belongs to a convex edge } \\ -2 \beta(x), & \text { if } x \text { belongs to a double convex edge } . \\ \beta(x)-\pi, & \text { if } x \text { belongs to a convex edge }\end{cases}
$$



Figure 22: The two different scenarios of intersections of the hyperplane at the three different types of edges appearing in a voxel manifold $\mathcal{M}_{\mathcal{V}}$ illustrated within the hyperplane $x+\mathcal{T}_{x} \partial_{1} \mathcal{M}_{\mathcal{V}}$. The dashed red line is the boundary of an $\delta$-ball centered at $x$. The grey shaded area belongs to $\mathcal{M}_{\mathcal{V}}$ and the red shaded area is the (geometric) normal cone $x+\mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}$. Note that the boundary of the normal cone is orthogonal to the boundary of $\mathcal{M}_{\mathcal{V}}$ The red arrow is the unit direction $-\nu \in \mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{1}{\nu}\right)$ scaled by $\epsilon$ and the turquoise line is the hyperplane defined by $-\nu$.

Proof. Let $\beta(x)$ be the minimum of $\varphi$ and $2 \pi-\varphi$ where $\phi$ is the opening angle within $\mathcal{M}_{\mathcal{V}}$ from Figure 22. This yields

$$
\begin{aligned}
\int_{\mathbb{S}\left(\mathcal{T}_{x} \partial_{1} \mathcal{M}_{\left.\frac{1}{\mathcal{V}}\right)}\right.} \alpha\left(\nu_{2}\right) \mathcal{H}_{1}\left(d \nu_{2}\right) & =\int_{0}^{2 \pi} \mathbb{1}_{-\sin (t) E_{1}-\cos (t) E_{2} \in\left(\mathcal{N}_{x} \mathcal{M}_{\mathcal{V}}\right)^{\circ} d t} \\
& = \begin{cases}\pi-\beta(x), & \text { if } x \text { on a convex edge } \\
-2 \beta(x), & \text { if } x \text { on a double convex edge }, \\
\beta(x)-\pi, & \text { if } x \text { on a convex edge }\end{cases}
\end{aligned}
$$

because the characteristic function is only 1 , if $-\nu$ belongs to the normal cone and hence the integral is equal to the opening angle of the normal cone (red shaded area in Fig. 22).

It remains to compute the angle $\beta(x)$ which is obtained by computing the angle between the intersection of the affine plane $x+\mathcal{T}_{x} \partial_{1} \mathcal{M} \frac{\perp}{\mathcal{V}}$ and the boundary of $\mathcal{M}_{\mathcal{V}}$ at $x$. We only treat the case of $x \in \mathcal{M}_{\mathcal{V}}$ lying on a convex edge. Double convex and concave edges follow analogously. Since the tangent space and $\mathbb{R}^{3}$ can be identified we assume w.l.o.g. that $x=0$ and the voxel is given by the set $\left\{y \in \mathbb{R}^{3}: y_{3} \leq 0, y_{2} \leq 0\right\}$. (note we extend the edge infinitely, which does not make a difference in this argument) Its boundary is given by the set $A \cup B=\left\{y \in \mathbb{R}^{3}: y_{3}=0, y_{2} \leq 0\right\} \cup\left\{y \in \mathbb{R}^{3}: y_{3} \leq 0, y_{2}=0\right\}$. The edge to which $x$ belongs is given by $E=\left\{y \in \mathbb{R}^{3}: y_{3}=0, y_{2}=0\right\}$. An orthonormal basis at $x$ with $U_{x}$ spanning $\mathcal{T}_{x} E$ is given in Proposition 15. The plane in which the unit circle $\mathbb{S}^{1}\left(\mathcal{T}_{x} E^{\perp}\right)$ lies is given by the linear span of $V_{x}, W_{x}$, which we denote by $\mathcal{F}=\left\{y \in \mathbb{R}^{3}: m_{1}(x) y_{1}+m_{2}(x) y_{2}+m_{3}(x) y_{3}=0\right\}$ for
some $m_{1}(x), m_{2}(x), m_{3}(x) \in \mathbb{R}$. The intersection $\mathcal{F} \cap A$ and $\mathcal{F} \cap B$ are given by

$$
\begin{align*}
& A \cap \mathcal{F}=\left\{y \in \mathbb{R}^{3}: m_{1}(x) y_{1}+m_{2}(x) y_{2}=0 \wedge y_{2} \leq 0 \wedge y_{3}=0\right\}  \tag{53}\\
& B \cap \mathcal{F}=\left\{y \in \mathbb{R}^{3}: m_{1}(x) y_{1}+m_{3}(x) y_{3}=0 \wedge y_{3} \leq 0 \wedge y_{2}=0\right\}
\end{align*}
$$

By construction $m_{1}(x) \neq 0$, since otherwise $\mathcal{F} \cap A=\mathcal{F} \cap B=E$ and hence $V, W$ cannot be both orthogonal to $E$ which contradicts the assumption that $U_{x}, V_{x}, W_{x}$ form an orthonormal basis for $\mathcal{T}_{x} \overline{\mathcal{M}_{\mathcal{V}}}$. Thus, if $m_{1}(x) \neq 0$ we have that $\left(m_{2}(x) / m_{1}(x),-1,0\right)$ and $\left(m_{3}(x) / m_{1}(x), 0,-1\right)$ are vectors in the intersection, which we can identify with tangent directions along $E$. Thus,

$$
\left.\begin{array}{rl}
\varphi & =\arccos \left(\frac{\frac{m_{2}(x) m_{3}(x)}{m_{1}^{2}(x)}}{\sqrt{\frac{m_{2}^{2}(x)}{m_{1}^{2}(x)}+1} \sqrt{\frac{m_{3}^{2}(x)}{m_{1}^{2}(x)}}+1}\right. \tag{54}
\end{array}\right) .
$$

The same formula holds true if $m_{1}(x)=0$ and $m_{2}(x) \neq 0$ and $m_{3}(x) \neq 0$.
Computation of $\int_{\partial_{2} M} \boldsymbol{\Lambda}_{x}\left(\bar{\nabla}_{U_{x}} U_{x}+\bar{\nabla}_{V_{x}} V_{x}, N_{x}\right) \mathcal{H}_{2}(d x)$ for Voxel Manifolds Assume w.l.o.g. that $\mathcal{T}_{x} \mathcal{M}_{\mathcal{V}}$ for $x \in \partial_{2} \mathcal{M}_{\mathcal{V}}$ is contained in the $x_{1}-x_{2}$-plane. An orthonormal frame is given by $U_{x}, V_{x}, W_{x}$ from (15) and by construction $N_{x}= \pm W_{x}$, where the sign depends on whether $W_{x}$ is inward or outward pointing. Using the coordinate representations $U_{x}=$ $\sum_{d=1}^{3} U_{d}(x) \partial_{d}$ and $V_{x}=\sum_{d=1}^{3} V_{d}(x) \partial_{d}$, linearity and product rule for the covariant derivative and formula (34), we obtain

$$
\begin{aligned}
\boldsymbol{\Lambda}_{x}\left(\bar{\nabla}_{U_{x}} U_{x}, N_{x}\right) & =N_{x}^{T} \boldsymbol{\Lambda}(x) U_{1}(x)\left(\begin{array}{c}
\partial_{1} U_{1}(x) \\
0 \\
0
\end{array}\right)+U_{1}^{2}(x) N_{x}^{T}\left(\begin{array}{l}
\Gamma_{111}(x) \\
\Gamma_{112}(x) \\
\Gamma_{113}(x)
\end{array}\right) \\
& =U_{1}^{2}(x) N_{x}^{T}\left(\begin{array}{l}
\Gamma_{111}(x) \\
\Gamma_{112}(x) \\
\Gamma_{113}(x)
\end{array}\right) .
\end{aligned}
$$

Here the second equality is due to the fact that $N_{x} \propto \boldsymbol{\Lambda}^{-1} E_{3}$, if represented as a vector. Similarly it holds that

$$
\begin{gathered}
\boldsymbol{\Lambda}_{x}\left(\bar{\nabla}_{V_{x}} V_{x}, N_{x}\right)=V_{1}^{2}(x) N_{x}^{T}\left(\begin{array}{l}
\Gamma_{111}(x) \\
\Gamma_{112}(x) \\
\Gamma_{113}(x)
\end{array}\right)+V_{2}^{2}(x) N_{x}^{T}\left(\begin{array}{l}
\Gamma_{221}(x) \\
\Gamma_{222}(x) \\
\Gamma_{223}(x)
\end{array}\right) \\
+V_{1}(x) V_{2}(x) N_{x}^{T}\left(\begin{array}{l}
\Gamma_{121}(x) \\
\Gamma_{122}(x) \\
\Gamma_{123}(x)
\end{array}\right)
\end{gathered}
$$

Summarizing this yields the following proposition about the trace of the shape operator along $\partial_{2} \mathcal{M}_{\mathcal{V}}$, i.e., $\bar{g}\left(\bar{\nabla}_{U_{x}} U_{x}+\bar{\nabla}_{V_{x}} V_{x}, N_{x}\right)$.

Proposition 9. Let $\mathcal{M}_{\mathcal{V}}$ be a voxel manifold and assume that $x \in \partial_{2} \mathcal{M}_{\mathcal{V}}$ such that $\mathcal{T}_{x} \partial_{2} \mathcal{M}_{\mathcal{V}}$ is spanned by $E_{k}, E_{l}$. Then

$$
\begin{aligned}
& \boldsymbol{\Lambda}_{x}\left(\bar{\nabla}_{U_{x}} U_{x}+\bar{\nabla}_{V_{x}} V_{x}, N_{x}\right) \\
& =\left(U_{k}^{2}(x)+V_{k}^{2}\right) \nu^{T}\left(\begin{array}{l}
\Gamma_{k k 1}(x) \\
\Gamma_{k k 2}(x) \\
\Gamma_{k k 3}(x)
\end{array}\right)+V_{l}^{2}(x) \nu^{T}\left(\begin{array}{c}
\Gamma_{l l 1}(x) \\
\Gamma_{l l 2}(x) \\
\Gamma_{l l 3}(x)
\end{array}\right) \\
& \quad+V_{k}(x) V_{l}(x) \nu^{T}\left(\begin{array}{l}
\Gamma_{k l 1}(x) \\
\Gamma_{k l 2}(x) \\
\Gamma_{k l 3}(x)
\end{array}\right)
\end{aligned}
$$

## D. 8 Proof of Theorem 3

We prove this result in more generality as the restriction to SuRFs is not necessary. The proof extends that of Tay [2007] which established the result for the trivial linear model. To do so assume that $Y_{1}, \ldots, Y_{N} \sim Y$ is an iid sample of Gaussian random fields on a voxel manifold $\mathcal{M}_{\mathcal{V}}$. We restrict here to voxel manifolds to avoid working in a chart, however, the proof easily generalizes. This sample we represent as the vector $\boldsymbol{Y}=\left(Y_{1}, \ldots, Y_{N}\right)^{T}$ and therefore the corresponding estimator generalizing (17) is, for $x \in \mathcal{M}_{\mathcal{V}}$,

$$
\begin{align*}
\hat{\Lambda}_{d d^{\prime}}(x)= & \frac{\operatorname{Cov}\left[\partial_{d} \boldsymbol{Y}(x), \partial_{d^{\prime}} \boldsymbol{Y}(x)\right]}{\operatorname{Var}[\boldsymbol{Y}(x)]} \\
= & \operatorname{Cov}\left[\partial_{d}\left(\frac{\boldsymbol{Y}(x)}{\sqrt{\operatorname{Var}[\boldsymbol{Y}(x)]}}\right), \partial_{d^{\prime}}\left(\frac{\boldsymbol{Y}(x)}{\sqrt{\operatorname{Var}[\boldsymbol{Y}(x)]}}\right)\right]  \tag{55}\\
& \quad-\frac{\operatorname{Cov}\left[\partial_{d} \boldsymbol{Y}(x), \boldsymbol{Y}(x)\right] \operatorname{Cov}\left[\boldsymbol{Y}(x), \partial_{d^{\prime}} \boldsymbol{Y}(x)\right]}{\operatorname{Var}[\boldsymbol{Y}(x)]^{2}} .
\end{align*}
$$

Here as in the main manuscript the operation $\operatorname{Var}[\cdot]$ and $\operatorname{Cov}[\cdot, \cdot]$ denote the sample variance and sample covariance respectively. We define the vector of normalized residuals to be $\mathbf{R}(x)=$ $\frac{\mathbf{H Y}(x)}{\|\mathbf{H} \boldsymbol{Y}(x)\|}, x \in \mathcal{M}_{\mathcal{V}}$, where $\|\cdot\|$ denotes the Euclidean norm and $\mathbf{H}=I_{N \times N}-\mathbb{1}^{T}$ with $\mathbb{1}^{T}=$ $(1, \ldots, 1) \in \mathbb{R}^{N}$, is the centering matrix. Note that $\mathbf{R}$ does not depend on the (unknown) variance $\operatorname{Var}[Y]$ as for the sample $\tilde{\boldsymbol{Y}}=\boldsymbol{Y} / \sqrt{\operatorname{Var}[Y]}$ we have that $\tilde{\mathbf{R}}(x)=\mathbf{R}(x)$ for all $x \in \mathcal{M}_{\mathcal{V}}$. Thus, using

$$
\nabla \mathbf{X}(x)=\left(\begin{array}{ccc}
\frac{\partial X_{1}}{\partial x_{1}}(x) & \ldots & \frac{\partial X_{1}}{\partial x_{D}}(x) \\
\vdots & \ddots & \vdots \\
\frac{\partial X_{D}}{\partial x_{1}}(x) & \ldots & \frac{\partial X_{D}}{\partial x_{D}}(x)
\end{array}\right)
$$

for $\mathbf{X}=\left(X_{1}, \ldots, X_{D}\right)^{T} \in C\left(\mathcal{M}_{\mathcal{V}}, \mathbb{R}^{D}\right)$, we can rewrite (55) in matrix terms as

$$
\hat{\mathbf{\Lambda}}=(\nabla \mathbf{R})^{T} \nabla \mathbf{R} .
$$

Recall that the true underlying $\boldsymbol{\Lambda}$ is given by

$$
\boldsymbol{\Lambda}=\mathbb{E}\left[\left(\nabla\left(\frac{Y-\mathbb{E}[Y]}{\sqrt{\operatorname{Var}[Y]}}\right)\right)^{T} \nabla\left(\frac{Y-\mathbb{E}[Y]}{\sqrt{\operatorname{Var}[Y]}}\right)\right]=\mathbb{E}\left[\left(\nabla \tilde{Y}^{T}\right) \nabla \tilde{Y}\right],
$$

where $\tilde{Y}=\frac{Y}{\sqrt{\operatorname{Var}[Y]}}$. Using this we obtain the following lemma which the proof of which follows the corresponding proof in Tay [2007].

Lemma 4. Under the assumption and notation described above, we have that

$$
\begin{equation*}
\mathbb{E}[\sqrt{\operatorname{det}(\hat{\boldsymbol{\Lambda}}(x))}]=\sqrt{\operatorname{det}(\boldsymbol{\Lambda}(x))} \tag{56}
\end{equation*}
$$

Proof. $\mathbf{H}$ is a projection matrix and so $\mathbf{H H}=\mathbf{H}$ and $\mathbf{H}^{T}=\mathbf{H}$. We can expand $\nabla \boldsymbol{R}$ in terms of $\tilde{\boldsymbol{Y}}$ as follows:

$$
\nabla \boldsymbol{R}=\frac{\nabla \mathbf{H} \tilde{\boldsymbol{Y}}}{\|\mathbf{H} \tilde{\boldsymbol{Y}}\|}-\frac{\mathbf{H} \tilde{\boldsymbol{Y}} \nabla\left((\mathbf{H} \tilde{\boldsymbol{Y}})^{T} \mathbf{H} \tilde{\boldsymbol{Y}}\right)}{2\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{3}}=\left(I_{N \times N}-\frac{\mathbf{H} \tilde{\boldsymbol{Y}}(\mathbf{H} \tilde{\boldsymbol{Y}})^{T}}{\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{2}}\right) \frac{\mathbf{H} \nabla \tilde{\boldsymbol{Y}}}{\|\mathbf{H} \tilde{\boldsymbol{Y}}\|}
$$

Now note that

$$
\mathbf{H}\left(I_{N \times N}-\frac{\mathbf{H} \tilde{\boldsymbol{Y}}(\mathbf{H} \tilde{\boldsymbol{Y}})^{T}}{\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{2}}\right) \mathbf{H}=\mathbf{H}-\frac{\mathbf{H} \tilde{\boldsymbol{Y}}(\mathbf{H} \tilde{\boldsymbol{Y}})^{T}}{\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{2}}
$$

and therefore it is idempotent with

$$
\operatorname{tr}\left(\mathbf{H}-\frac{\mathbf{H} \tilde{\boldsymbol{Y}}(\mathbf{H} \tilde{\boldsymbol{Y}})^{T}}{\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{2}}\right)=\operatorname{tr}(\mathbf{H})-1
$$

Thus, applying Cochran's Theorem [Mardia et al., 1979, Theorem 3.4.4], it follows that for all $x \in S$,

$$
(\nabla \boldsymbol{R}(x))^{T} \nabla \boldsymbol{R}(x) \mid \tilde{\boldsymbol{Y}}(x) \sim \text { Wish }_{D}\left(\boldsymbol{\Lambda}\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{-2}, \operatorname{tr}(\mathbf{H})-1\right)
$$

As such by [Mardia et al., 1979, Corollary 3.4.1.2]

$$
\begin{aligned}
\boldsymbol{W}(x) & \sim\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{2} \boldsymbol{\Lambda}(x)^{-1 / 2}(\nabla \boldsymbol{R}(x))^{T} \nabla \boldsymbol{R}(x) \boldsymbol{\Lambda}(x)^{-1 / 2} \mid \tilde{\boldsymbol{Y}}(x) \\
& \sim \text { Wish }_{D}\left(I_{N \times N}, \operatorname{tr}(\mathbf{H})-1\right)
\end{aligned}
$$

Taking determinants and rearranging it follows that, unconditionally,

$$
\sqrt{\operatorname{det}\left((\nabla \boldsymbol{R}(x))^{T} \nabla \boldsymbol{R}(x)\right)} \sim \sqrt{\operatorname{det}(\boldsymbol{\Lambda}(x)) \operatorname{det}(\boldsymbol{W}(x)) V(x)^{-D}}
$$

where $\boldsymbol{W}(x)$ is independent of $V(x)=\|\mathbf{H} \tilde{\boldsymbol{Y}}\|^{2} \sim \chi_{\operatorname{tr}(\mathbf{H})}^{2}$. From the independence, the formula for the moments of $\chi^{2}$-distributions and Theorem 3.4.8 from Mardia et al. [1979] we obtain

$$
\begin{aligned}
\mathbb{E}\left[\sqrt{\operatorname{det}\left((\nabla \boldsymbol{R}(x))^{T} \nabla \boldsymbol{R}(x)\right)}\right] & =\sqrt{\operatorname{det}(\boldsymbol{\Lambda}(x))} \mathbb{E}[\sqrt{\operatorname{det}(\boldsymbol{W}(x))}] \\
\mathbb{E}\left[V(x)^{-\frac{D}{2}}\right] & =\operatorname{det}(\boldsymbol{\Lambda}(x))^{1 / 2}
\end{aligned}
$$

Remark 12. Lemma 4 also holds for estimates of $\boldsymbol{\Lambda}$ in a linear model, compare Tay [2007]. This follows by a similar proof as, in that setting, the residuals $\mathbf{R}$ which are used to calculate $\hat{\boldsymbol{\Lambda}}$ are obtained by $\mathbf{R}=\mathbf{P Y}$, where $\mathbf{P}$ is idempotent.

Proof of Theorem 3. In what follows we establish the results for $\hat{\mathcal{L}}_{D}^{(r)}$. The proof for $\hat{\mathcal{L}}_{D-1}^{(r)}$ is identical as each $\boldsymbol{\Lambda}^{I}$ is a $(D-1) \times(D-1)$-submatrix of $\boldsymbol{\Lambda}$.

Using Lemma 4 and the approximation of integrals of continuous functions by Riemann sums yields

$$
\begin{aligned}
\lim _{r \rightarrow \infty} \mathbb{E}\left[\hat{\mathcal{L}}_{D}^{(r)}\right] & =\lim _{r \rightarrow \infty} \mathbb{E}\left[\sum_{x \in \mathcal{M}_{\nu}^{(r)}} \sqrt{\operatorname{det}(\hat{\boldsymbol{\Lambda}}(x))} \prod_{d=1}^{D} \frac{\delta_{d}}{r+1}\right] \\
& =\lim _{r \rightarrow \infty} \sum_{x \in \mathcal{M}_{\nu}^{(r)}} \mathbb{E}[\sqrt{\operatorname{det}(\hat{\boldsymbol{\Lambda}}(x))}] \prod_{d=1}^{D} \frac{\delta_{d}}{r+1} \\
& =\lim _{r \rightarrow \infty} \sum_{x \in \mathcal{M}_{\nu}^{(r)}} \sqrt{\operatorname{det}(\boldsymbol{\Lambda}(x))} \prod_{d=1}^{D} \frac{\delta_{d}}{r+1}=\mathcal{L}_{D}
\end{aligned}
$$

Similarly, applying Fubini's theorem which is applicable as $\mathcal{M}_{\mathcal{V}}$ is compact, we obtain that

$$
\begin{aligned}
\mathbb{E}\left[\lim _{r \rightarrow \infty} \hat{\mathcal{L}}_{D}^{(r)}\right] & =\mathbb{E}\left[\int_{\mathcal{M}_{\mathcal{V}}} \sqrt{\operatorname{det}(\hat{\boldsymbol{\Lambda}}(x))} \prod_{d=1}^{D} d x_{d}\right] \\
& =\int_{\mathcal{M}_{\mathcal{V}}} \mathbb{E}[\sqrt{\operatorname{det}(\hat{\boldsymbol{\Lambda}}(x))}] \prod_{d=1}^{D} d x_{d} \\
& =\int_{\mathcal{M}_{\mathcal{V}}} \sqrt{\operatorname{det}(\boldsymbol{\Lambda}(x))} \prod_{d=1}^{D} d x_{d}=\mathcal{L}_{D}
\end{aligned}
$$

## D. 9 Proof of Theorem 4

Proof. By the assumptions and Proposition 2 the SuRF has almost surely $L^{2}$-Hölder continuous paths and therefore the assumptions of Lemma 11 from Telschow and Schwartzman [2022] are satisfied. This means that $\hat{\boldsymbol{\Lambda}}$ converges uniformly almost surely to $\boldsymbol{\Lambda}$ over all $x \in \overline{\mathcal{M}_{\mathcal{V}}}$. Since the Riemann sum converges to the integral the $\lim _{r \rightarrow \infty} \lim _{N \rightarrow \infty}$ statement follows immediately. On the other hand the $\lim _{N \rightarrow \infty} \lim _{r \rightarrow \infty}$ statement is a special case of Theorem 3 from Telschow et al. [2023] since for $\mathcal{L}_{2}$ and $\mathcal{L}_{3}$ from their condition (R) only the uniform almost sure convergence of $\hat{\boldsymbol{\Lambda}}$ to $\boldsymbol{\Lambda}$ is required.


[^0]:    ${ }^{1}$ We favor the term atoms over data objects because this nomenclature better separates the two concepts.

[^1]:    ${ }^{2}$ Stating the GKF requires the SuRFs to be defined on a $D$-dimensional, compact manifold without boundary $\overline{\mathcal{M}_{\nu}} \supset \mathcal{M}_{\nu}$

[^2]:    ${ }^{3}$ The two-sided hypothesis can be treated similarly.

[^3]:    ${ }^{4}$ Our arguments carry over to tests controlling the FWER in the strong sense wrt the smoothed signal, for example, permutation tests in fMRI as they are typically applied to smoothed data.

